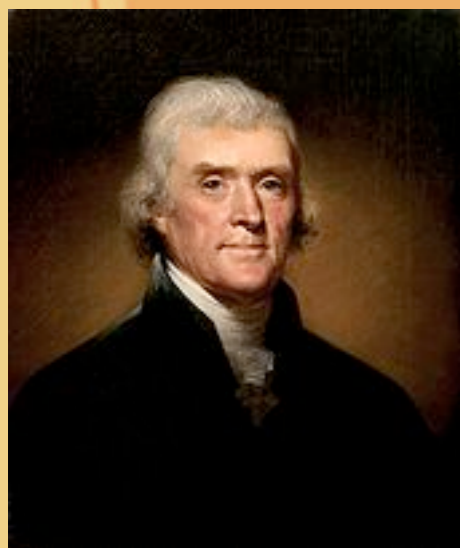




Modeling QCD for Hadron Physics



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Topics

- ◆ Overview of DSE modeling--mainly soft scale
- ◆ Masses, decays, form factors--mainly mesons
- ◆ Nucleon form factors
- ◆ Hard scale: DIS valence $u(x)$ in pion, kaon; qQ mesons
- ◆ Is $\langle \bar{q}q \rangle_\mu^0$ really an in-hadron condensate?



Lattice-QCD and DSE-based modeling

- Lattice: $\langle \mathcal{O} \rangle = \int D\bar{q}qG \mathcal{O}(\bar{q}, q, G) e^{-\mathcal{S}[\bar{q}, q, G]}$
 - Euclidean metric, x-space, Monte-Carlo
 - Issues: lattice spacing and vol, sea and valence m_q , fermion Det
 - **Large time limit** \Rightarrow nearest hadronic mass pole
- EOMs (DSEs): $0 = \int D\bar{q}qG \frac{\delta}{\delta q(x)} e^{-\mathcal{S}[\bar{q}, q, G] + (\bar{\eta}, q) + (\bar{q}, \eta) + (J, G)}$
 - Euclidean metric, p-space, continuum integral eqns
 - Issues: truncation and phenomenology – not full QCD
 - **Analytic contin.** \Rightarrow nearest hadronic mass pole
 - Can be quick to identify systematics, mechanisms, ...



DSE-based modeling of Hadron Physics

- Soft physics: truncate DSEs to min: 2-pt, 3-pt fns
- Should be **relativistically covariant**—convenient for decays, Form Factors, etc
 - **No boosts needed on wavefns of recoiling bound st.**
 - **∞ d.o.f. \rightarrow few quasi-particle effective d.o.f.**
- Do not make a 3-dimensional reduction
- Preserve 1-loop QCD renorm group behavior in UV
- Preserve global symmetries, conserved em currents, etc
- Preserve PCAC \Rightarrow Goldstone's Thm
- Can't preserve local color gauge covariance—just choose Landau gauge [RG fixed pt]
- Parameterize the deep infrared (large distance) QCD coupling





Constraints on Modeling

- Preserve vector WTI, and **axial vector WTI**

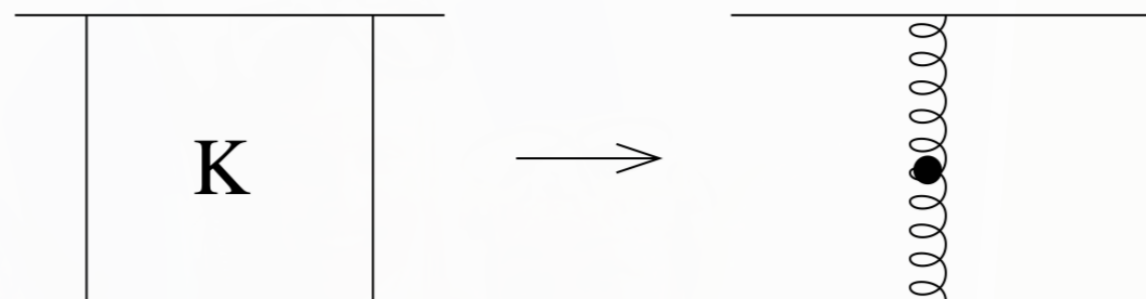
E.g.

$$-iP_\mu \Gamma_{5\mu}(k; P) = S^{-1}(k_+) \gamma_5 \frac{\tau}{2} + \gamma_5 \frac{\tau}{2} S^{-1}(k_-) - 2m_q(\mu) \Gamma_5(k; P)$$

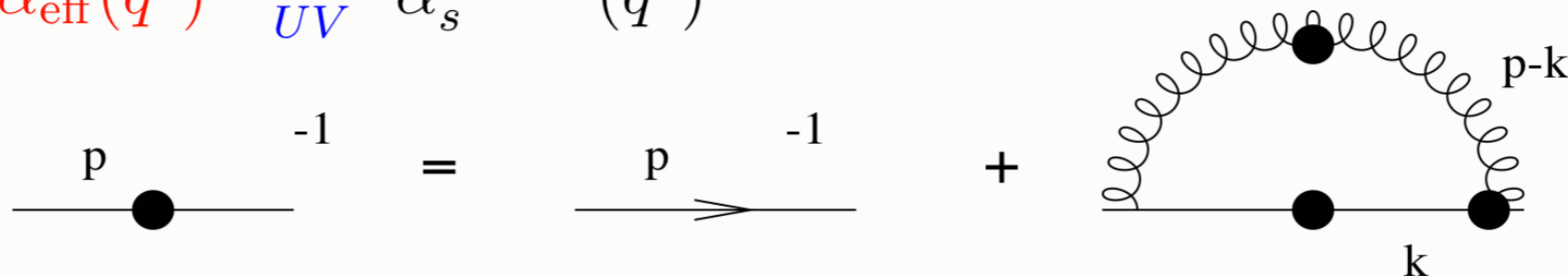
- \Rightarrow kernels of DSE_q and K_{BSE} are related
- Ladder-rainbow is the simplest implementation
- **Goldstone Theorem preserved**, ps octet masses good, indep of model details
- **DCSB** $\Rightarrow \pi$: $\Gamma_\pi^0(p^2) = \frac{i\gamma_5}{f_\pi^0} \left[\frac{1}{4} \text{tr} S_0^{-1}(p^2) \right] + \dots$
- Here, 1-body and 2-body systems are the same



Ladder-Rainbow Model



- $K_{\text{BSE}} \rightarrow -\gamma_\mu \frac{\lambda^a}{2} 4\pi\alpha_{\text{eff}}(q^2) D_{\mu\nu}^{\text{free}}(q) \gamma_\nu \frac{\lambda^a}{2}$
- $\alpha_{\text{eff}}(q^2) \xrightarrow{\overline{IR}} \langle \bar{q}q \rangle_{\mu=1} \text{ GeV} = -(240\text{MeV})^3$, incl vertex dressing
- $\alpha_{\text{eff}}(q^2) \xrightarrow{\overline{UV}} \alpha_s^{1\text{-loop}}(q^2)$



- P. Maris & P.C. Tandy, PRC60, 055214 (1999)
 M_ρ, M_ϕ, M_{K^*} good to 5%, f_ρ, f_ϕ, f_{K^*} good to 10%



Summary of light meson results

$m_{u=d} = 5.5 \text{ MeV}$, $m_s = 125 \text{ MeV}$ at $\mu = 1 \text{ GeV}$

Pseudoscalar (PM, Roberts, PRC56, 3369)

	expt.	calc.
$-\langle \bar{q}q \rangle_\mu^0$	$(0.236 \text{ GeV})^3$	$(0.241^\dagger)^3$
m_π	0.1385 GeV	0.138 [†]
f_π	0.0924 GeV	0.093 [†]
m_K	0.496 GeV	0.497 [†]
f_K	0.113 GeV	0.109

Charge radii (PM, Tandy, PRC62, 055204)

r_π^2	0.44 fm ²	0.45
$r_{K^+}^2$	0.34 fm ²	0.38
$r_{K^0}^2$	-0.054 fm ²	-0.086

$\gamma\pi\gamma$ transition (PM, Tandy, PRC65, 045211)

$g_{\pi\gamma\gamma}$	0.50	0.50
$r_{\pi\gamma\gamma}^2$	0.42 fm ²	0.41

Weak K_{l3} decay (PM, Ji, PRD64, 014032)

$\lambda_+(e3)$	0.028	0.027
$\Gamma(K_{e3})$	$7.6 \cdot 10^6 \text{ s}^{-1}$	7.38
$\Gamma(K_{\mu3})$	$5.2 \cdot 10^6 \text{ s}^{-1}$	4.90

Vector mesons (PM, Tandy, PRC60, 055214)

$m_{\rho/\omega}$	0.770 GeV	0.742
$f_{\rho/\omega}$	0.216 GeV	0.207
m_{K^*}	0.892 GeV	0.936
f_{K^*}	0.225 GeV	0.241
m_ϕ	1.020 GeV	1.072
f_ϕ	0.236 GeV	0.259

Strong decay (Jarecke, PM, Tandy, PRC67, 035202)

$g_{\rho\pi\pi}$	6.02	5.4
$g_{\phi KK}$	4.64	4.3
$g_{K^*K\pi}$	4.60	4.1

Radiative decay (PM, nucl-th/0112022)

$g_{\rho\pi\gamma}/m_\rho$	0.74	0.69
$g_{\omega\pi\gamma}/m_\omega$	2.31	2.07
$(g_{K^*K\gamma}/m_{K^*})^+$	0.83	0.99
$(g_{K^*K\gamma}/m_{K^*})^0$	1.28	1.19

Scattering length (PM, Cotanch, PRD66, 116010)

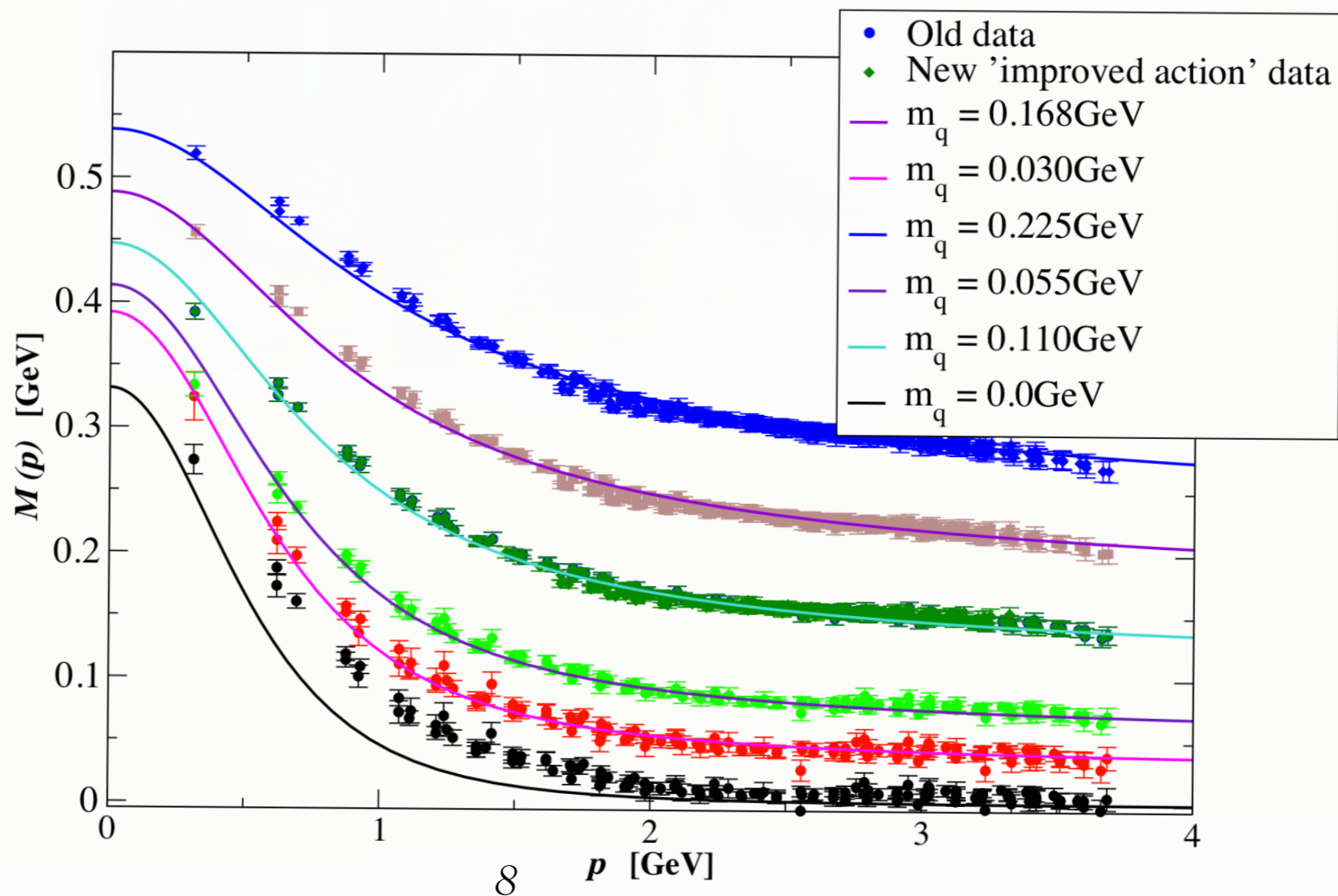
a_0^0	0.220	0.170
a_0^2	0.044	0.045
a_1^1	0.038	0.036

bsampl



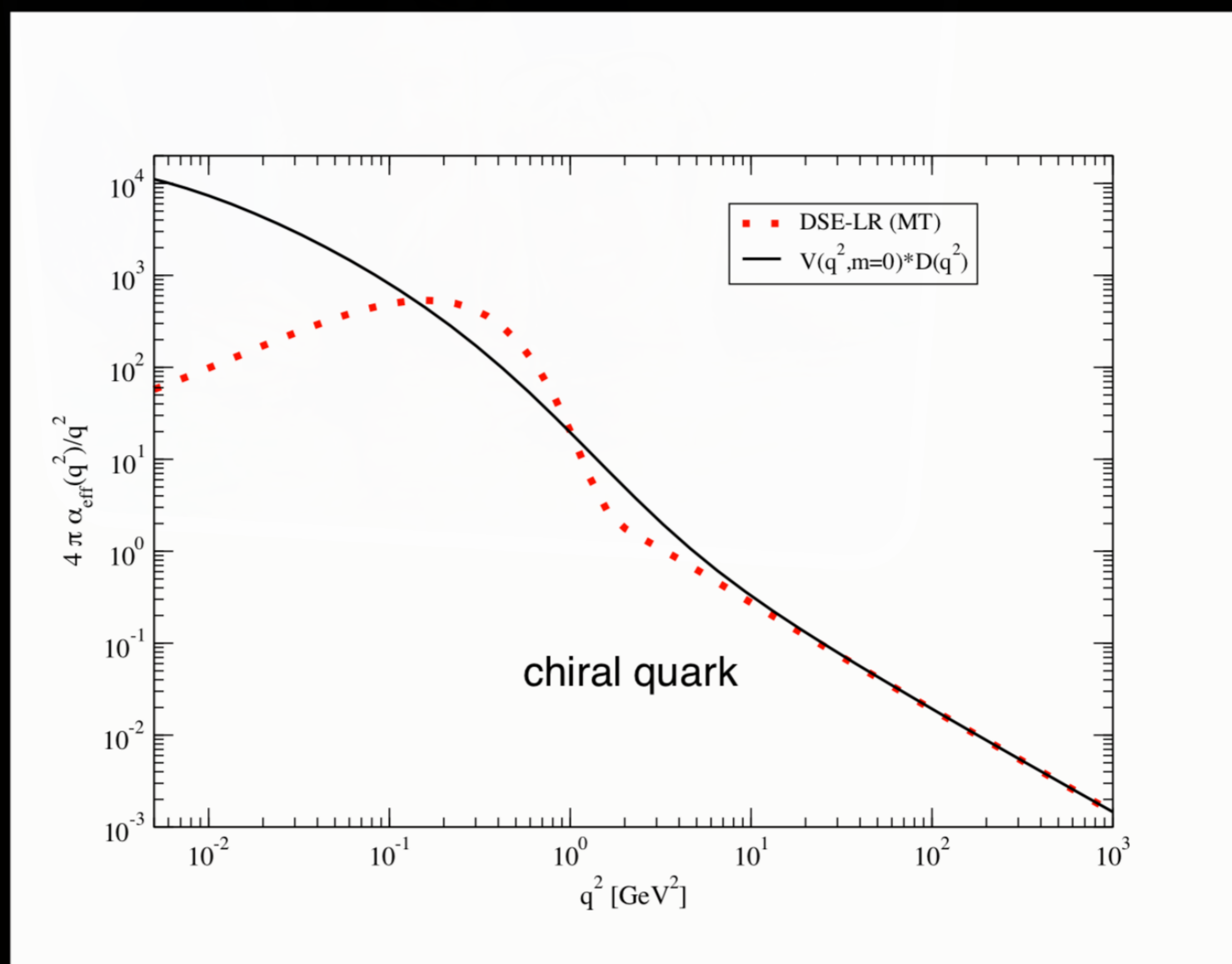
Qu-lattice $S(p)$, $D(q)$ mapped to a DSE kernel

$$S(p) = Z(p) [i \not{p} + M(p)]^{-1}$$





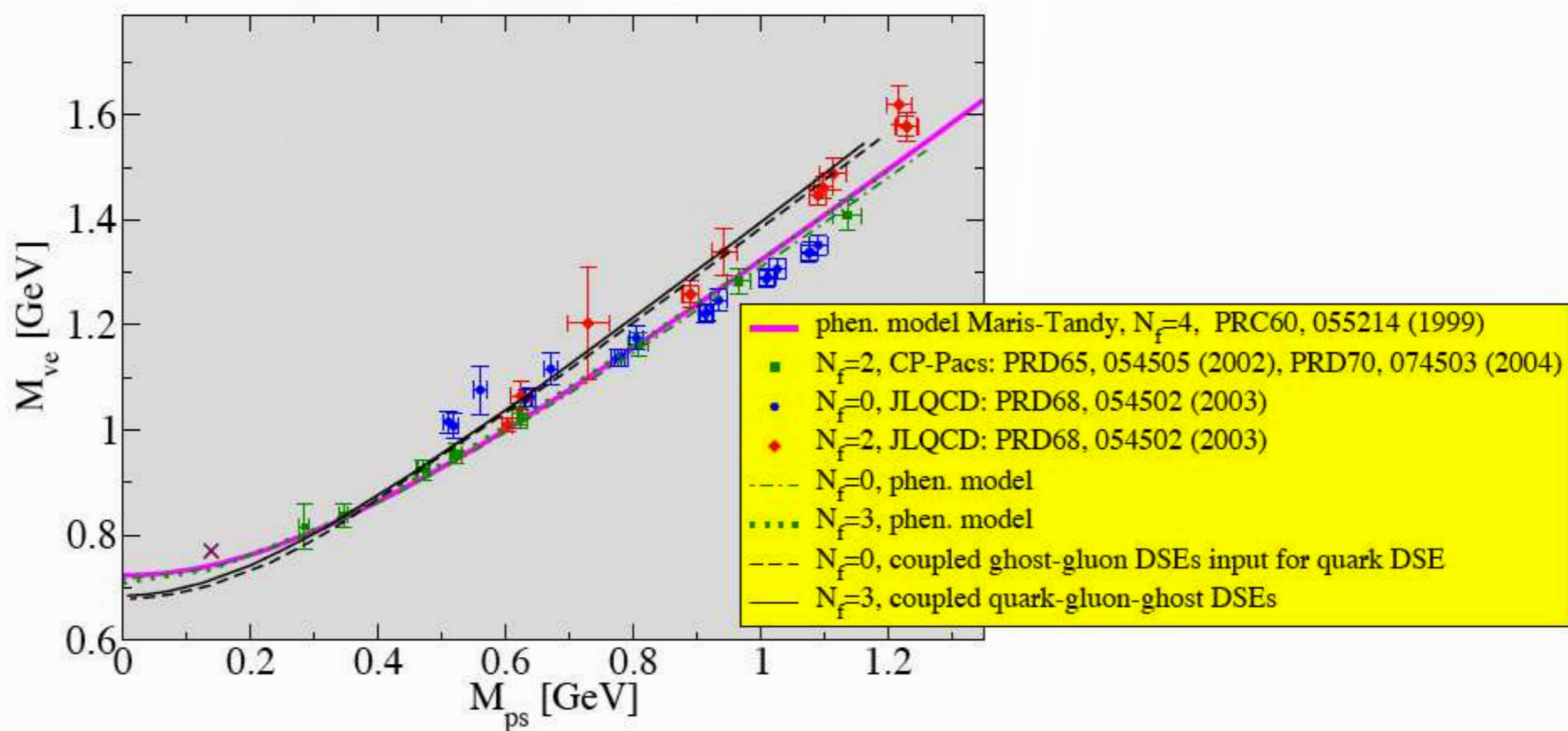
Quenched lattice $\Rightarrow m_q$ Depn of DSE Kernel



Bhagwat, Pichowsky, Roberts, Tandy, PRC68, 015203 (2003)



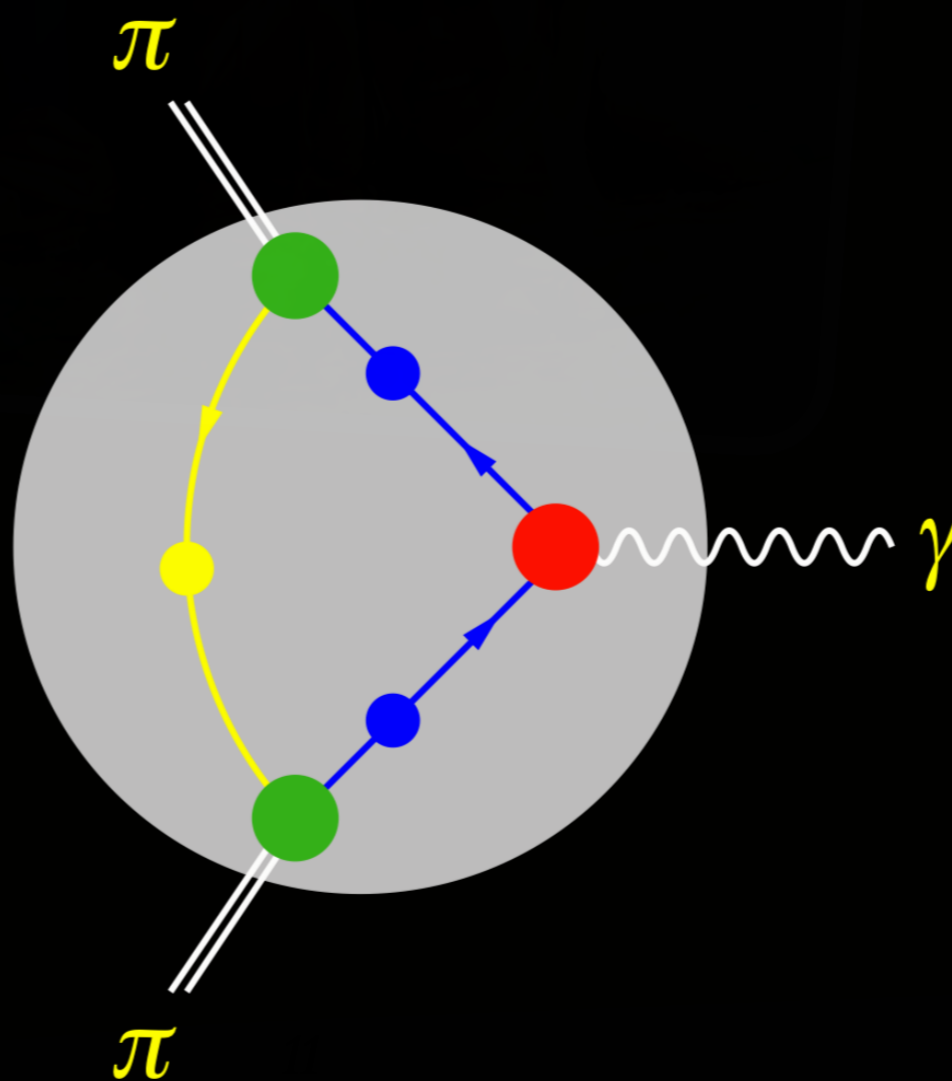
DSE and Lattice results for M_V and M_{ps}





Pion electromagnetic form factor

$$\Lambda_\mu = (P' + P)_\mu F_\pi(Q^2) = N_c \int \frac{d^4 q}{(2\pi)^4} \text{Tr} [\bar{\Gamma}^\pi S i\Gamma_\mu S \Gamma^\pi S]$$



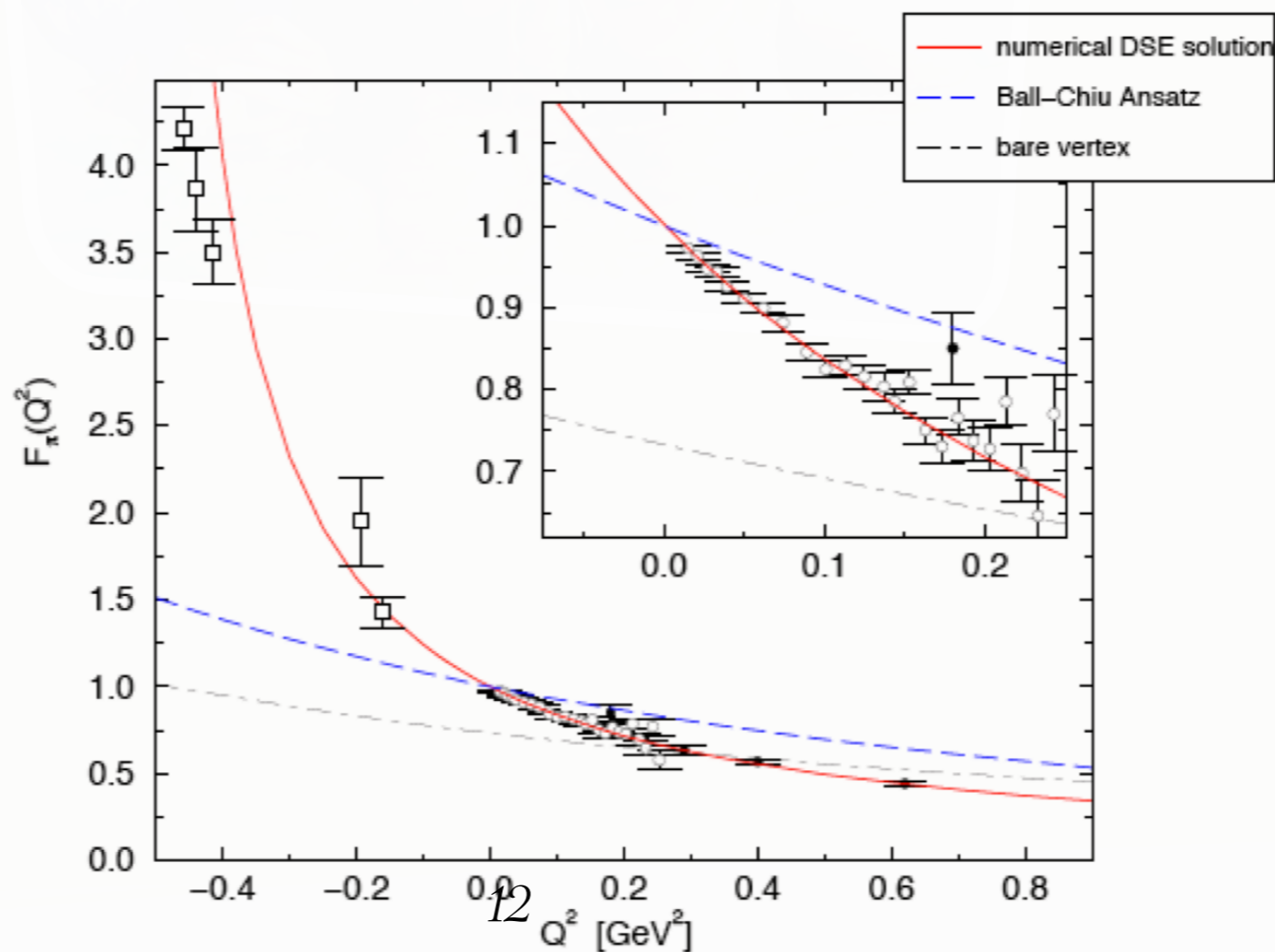


Pion $F(Q^2)$: Low Q^2

(P Maris & PCT, PRC 61, 045202 (2000))

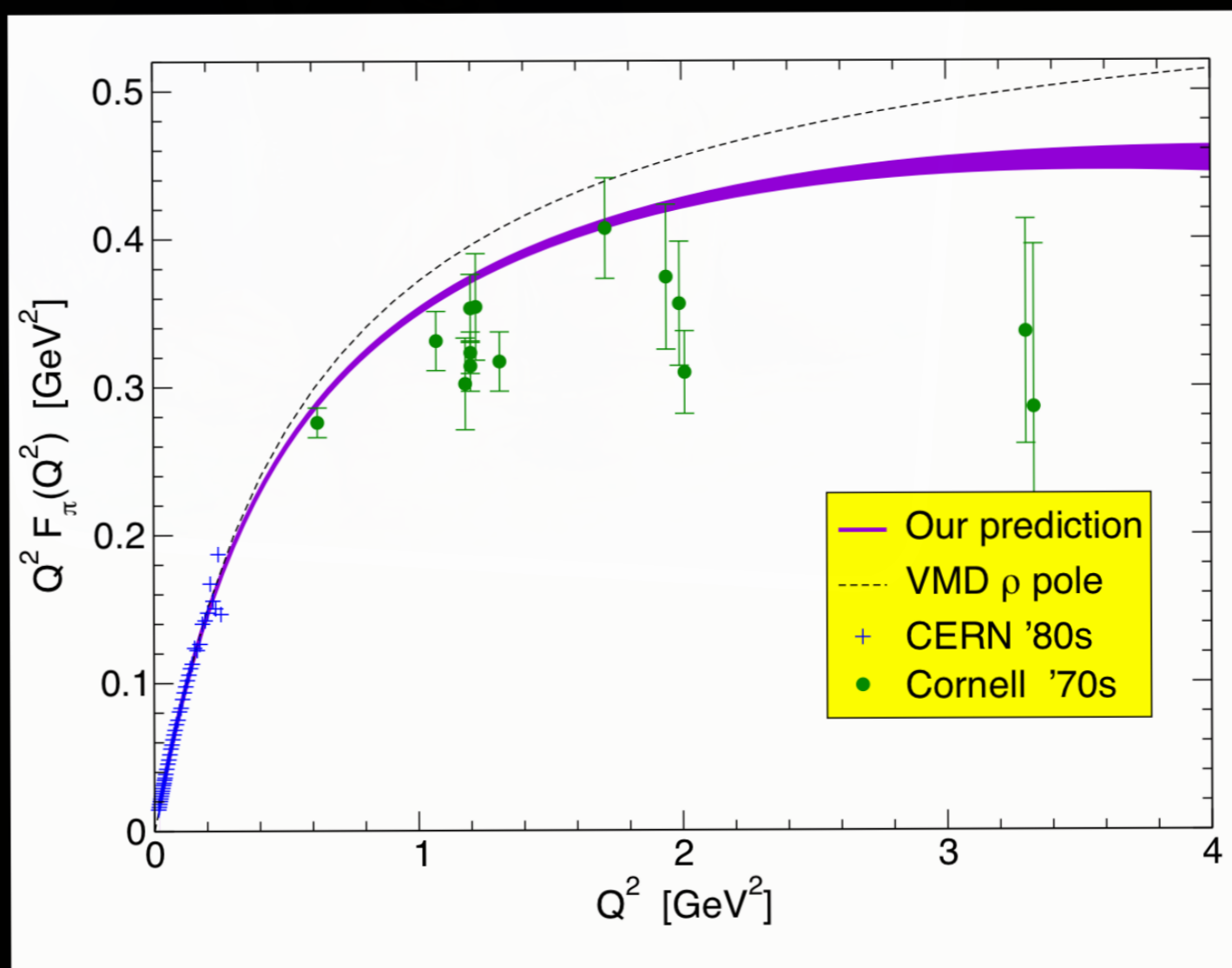
(P. Maris & PCT, PRC 62, 0555204 (2000))

$$r_{\pi}^{\text{DSE}} = 0.68 \text{ fm} \quad r_{\pi}^{\text{expt}} = 0.663 \pm .006 \text{ fm}$$





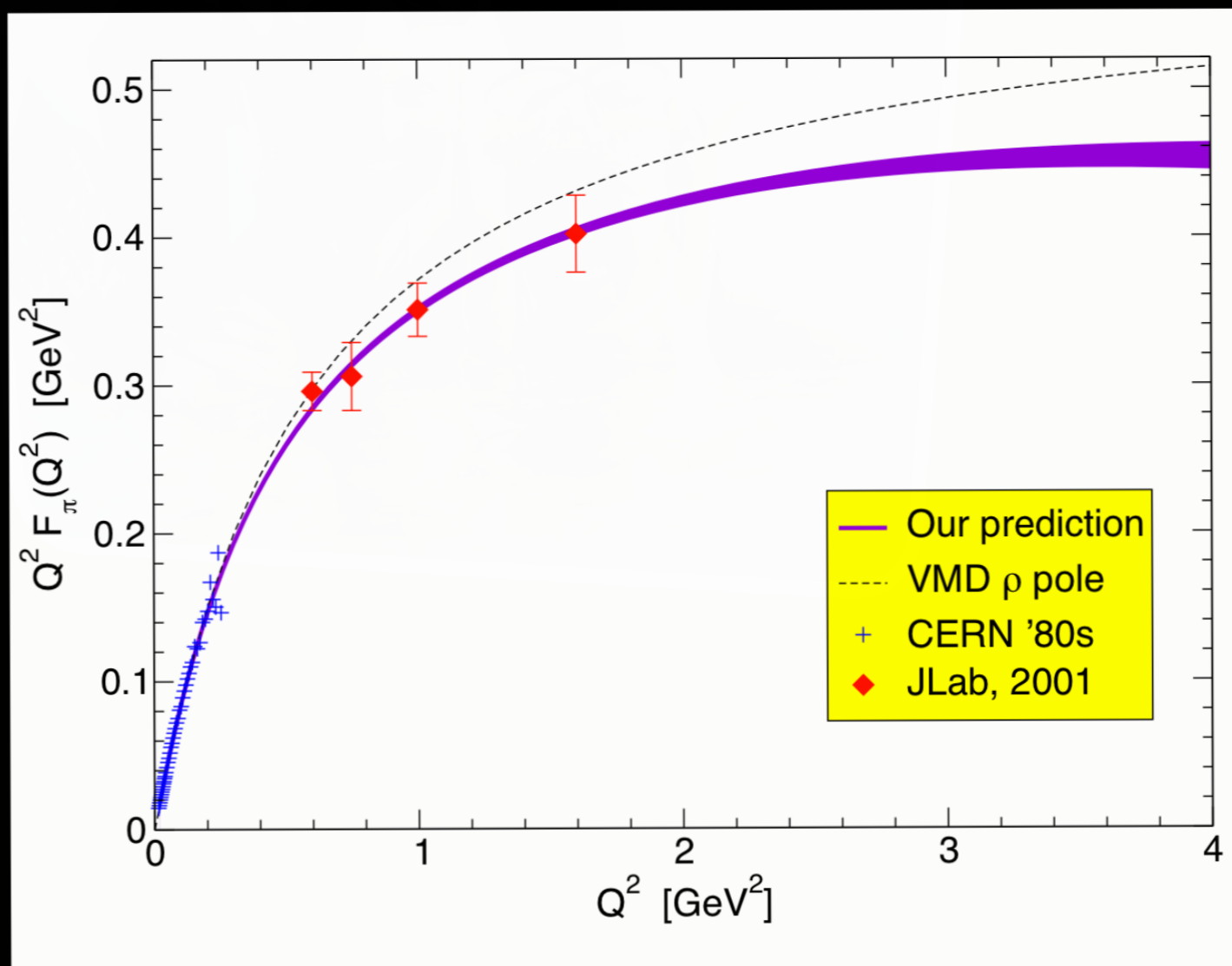
Pion electromagnetic form factor



PM and Tandy, PRC62,055204 (2000) [nucl-th/0005015]



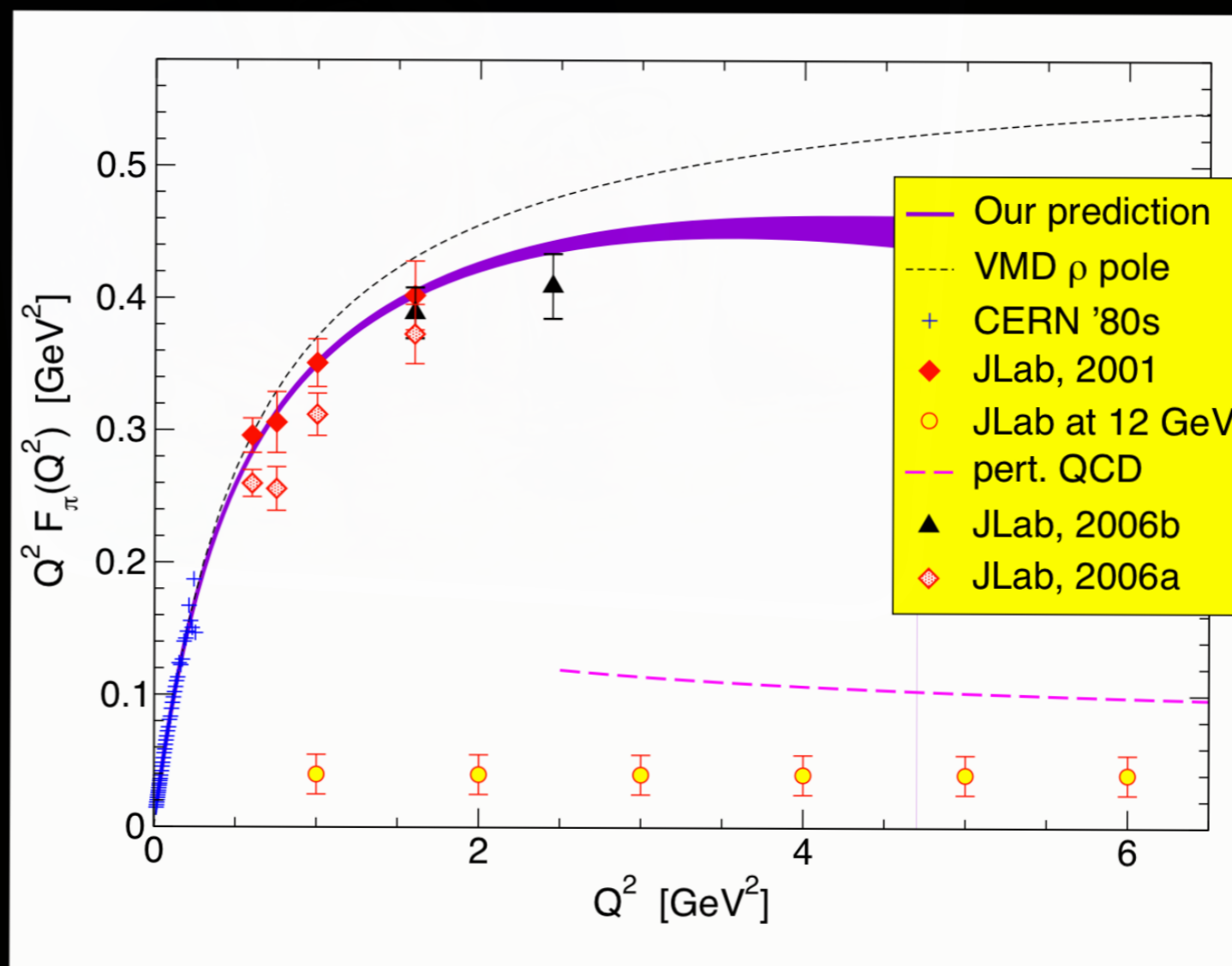
Pion electromagnetic form factor



JLab data from Volmer *et al*, PRL86, 1713 (2001) [nucl-ex/0010009]
PM and Tandy, PRC62,055204 (2000) [nucl-th/0005015]



Pion electromagnetic form factor

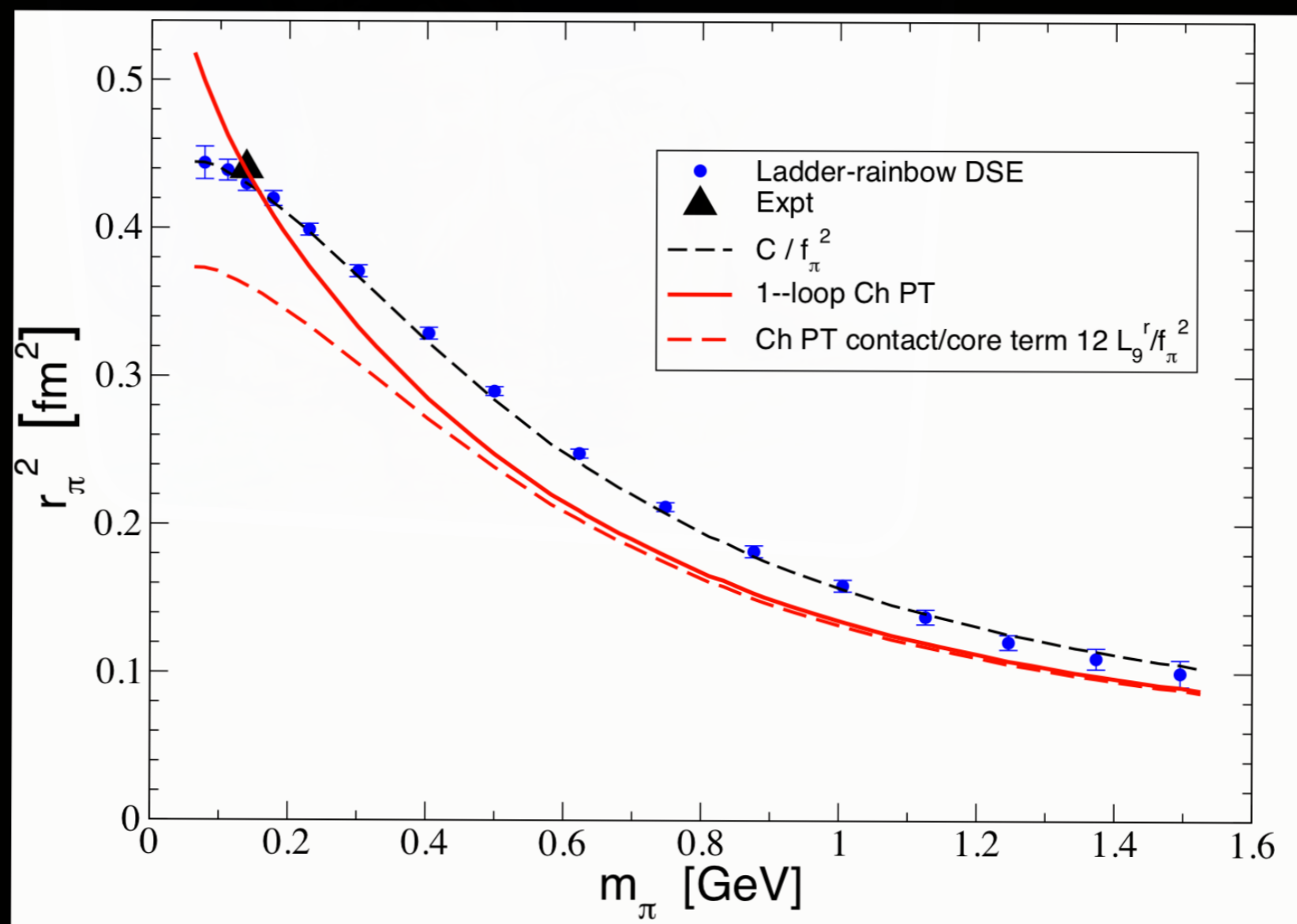


PM and Tandy, PRC62,055204 (2000) [nucl-th/0005015]

2006a: V. Tadevosyan *et al*, [nucl-ex/0607007], 2006b: T. Horn *et al*, [nucl-ex/0607005]



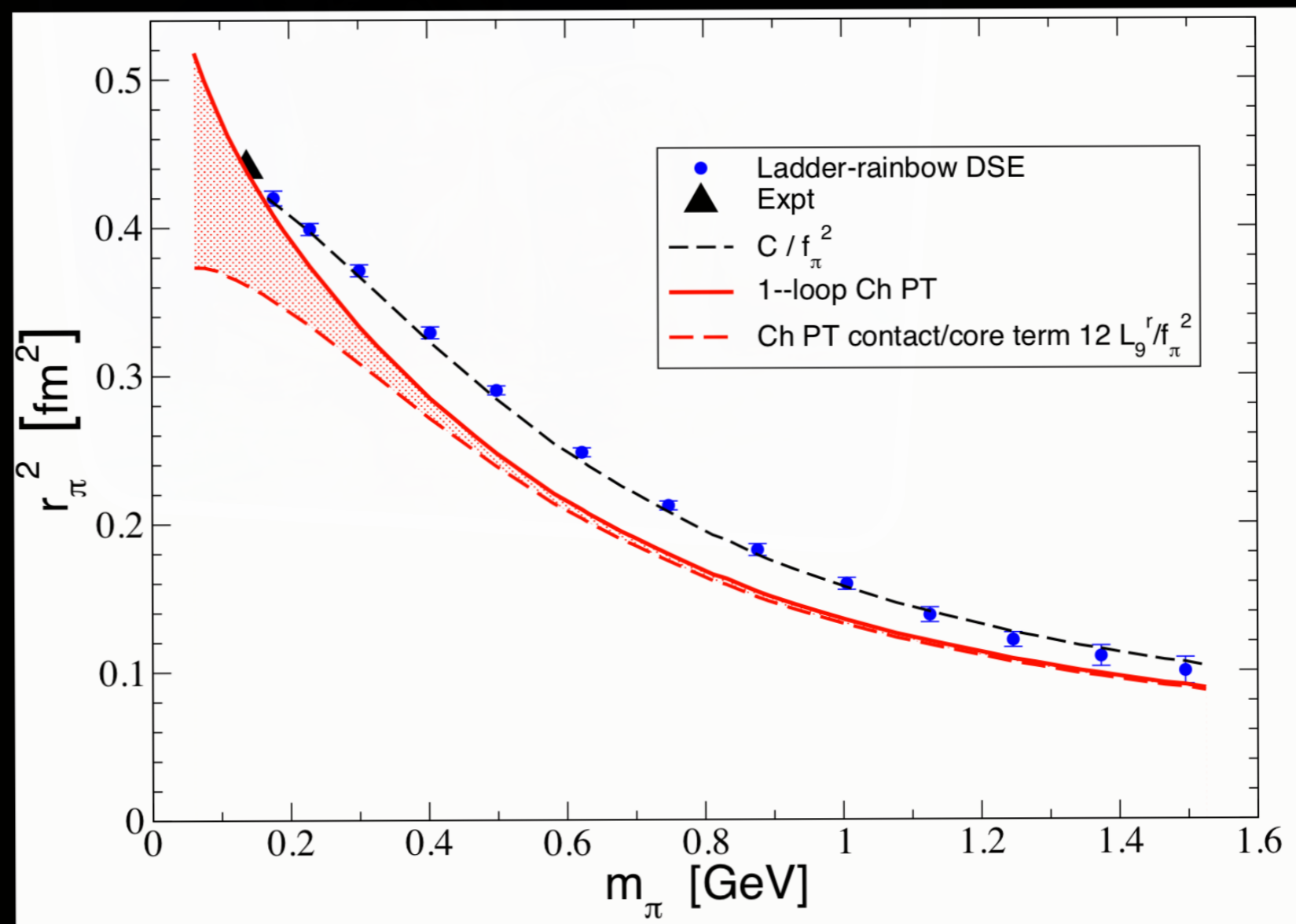
1-loop chiral correction to r_π vs m_π



P. Maris and PCT, in preparation



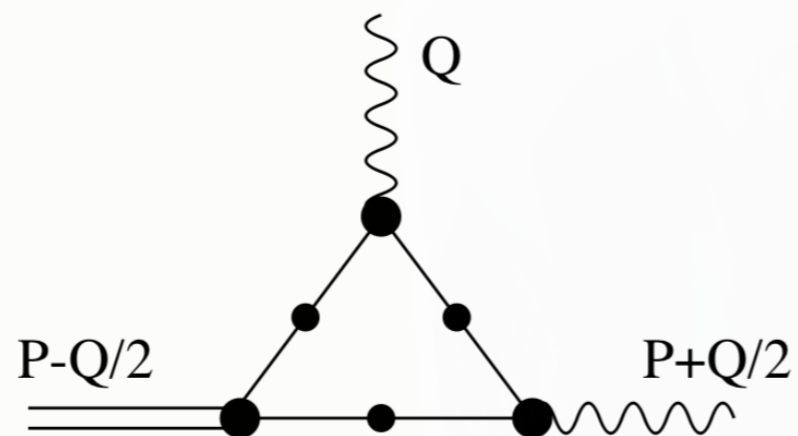
1-loop chiral correction to r_π vs m_π



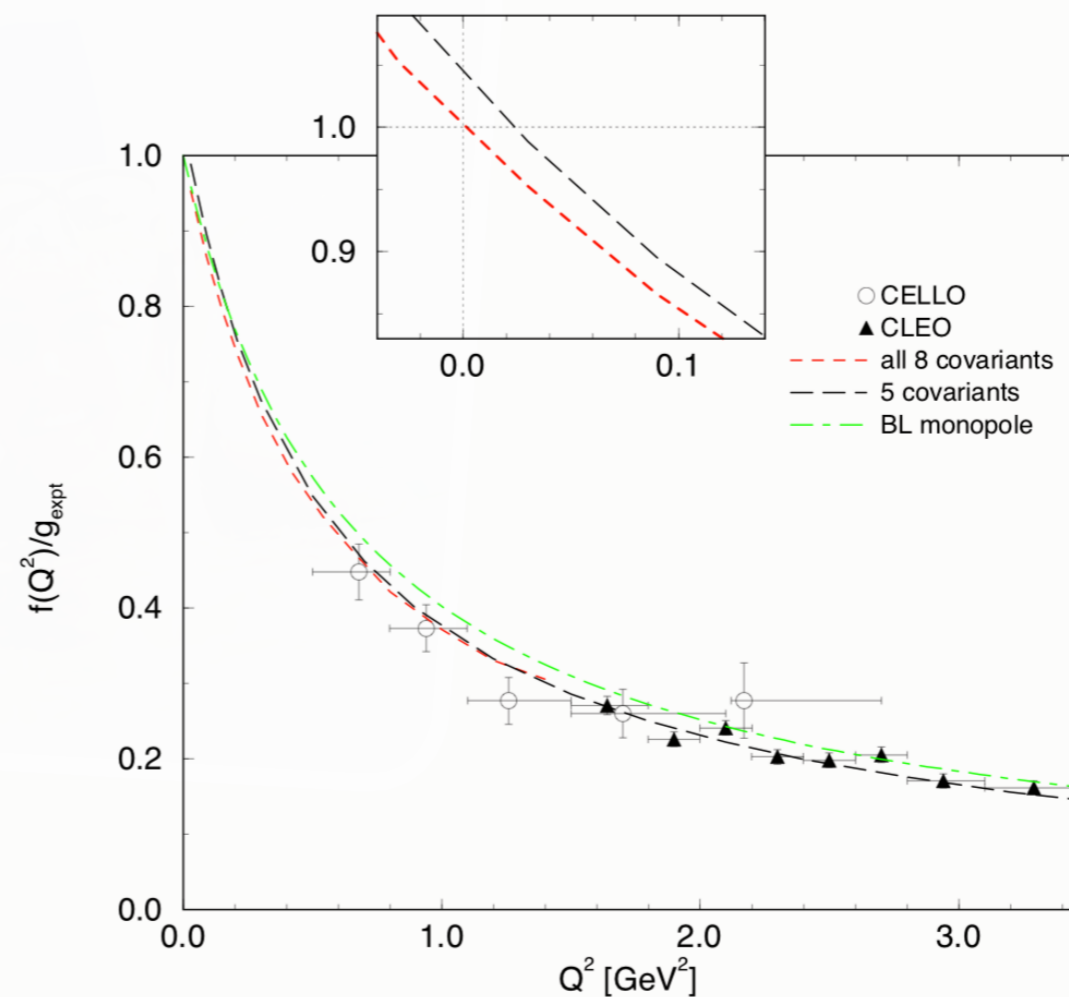
P. Maris and PCT, in preparation



$\gamma^* \pi^0 \rightarrow \gamma$ Transition Form Factor



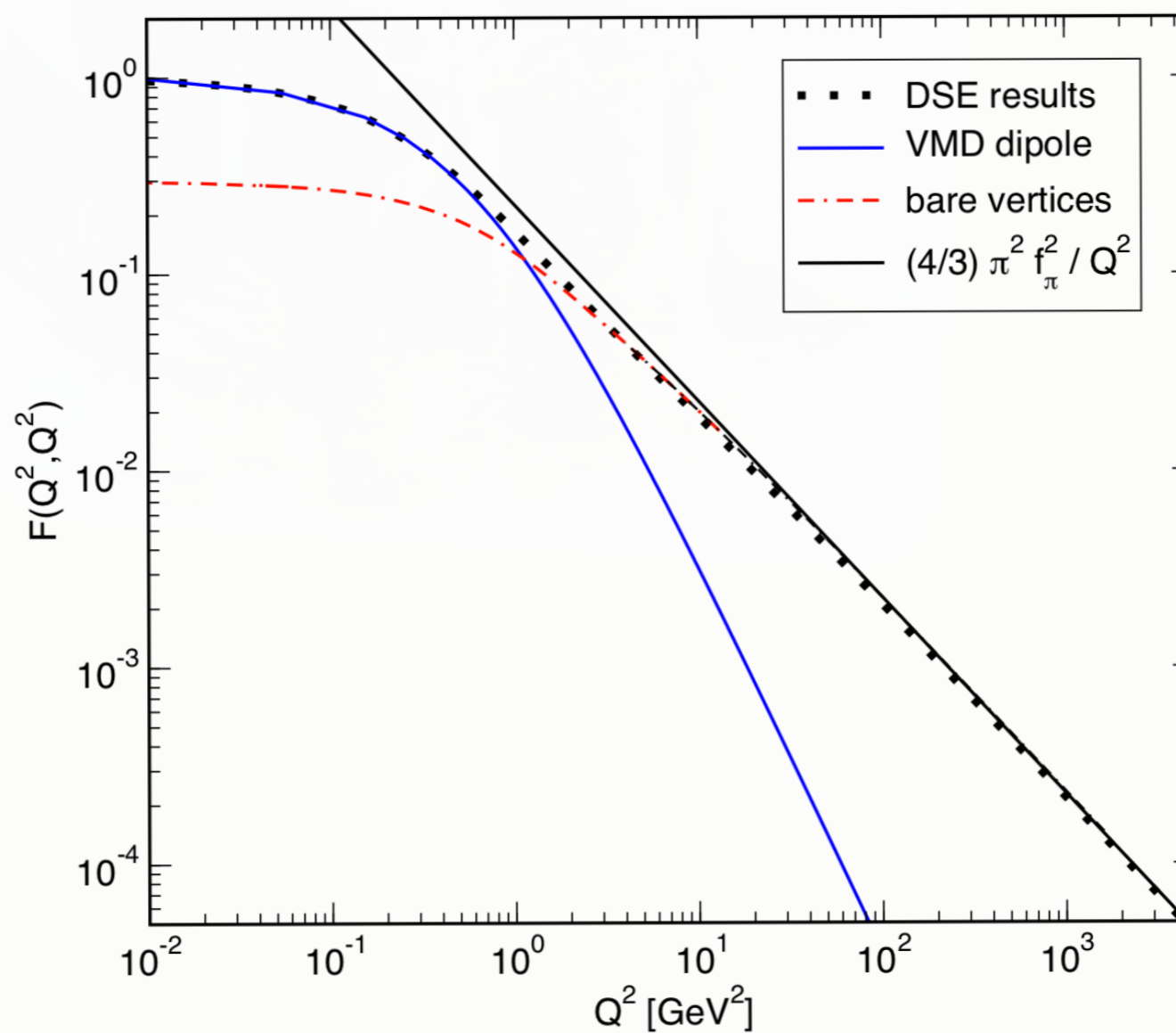
- Abelian axial anomaly + π pole
in $\Gamma_{5\mu} \Rightarrow G(0, 0)$
- Chiral limit $G(0, 0) = \frac{1}{2}$
 $\Rightarrow \Gamma_{\pi\gamma\gamma}$ to 2%



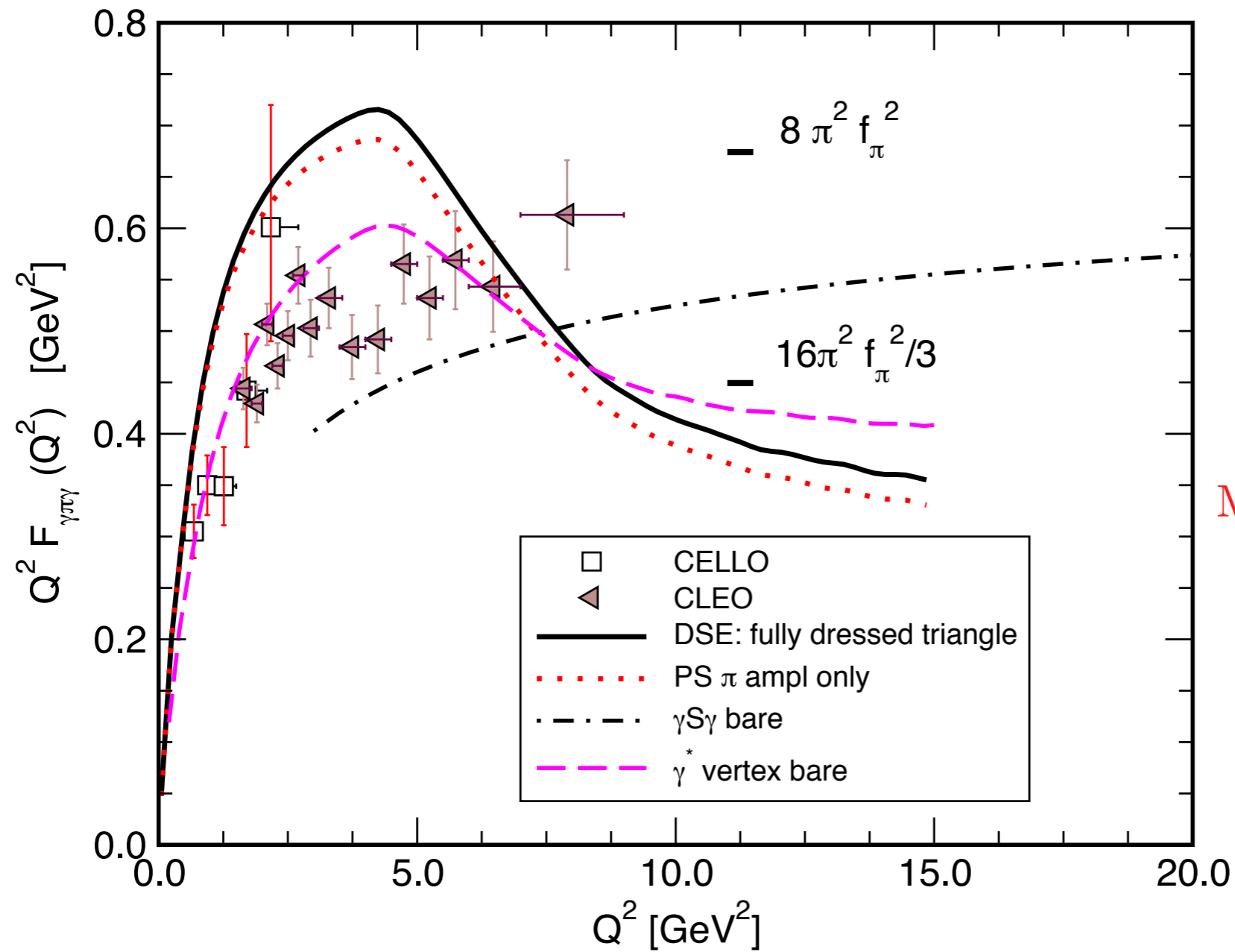


$\gamma^* \pi \gamma^*$ *Asymptotic Limit*

Lepage and Brodsky, PRD22, 2157 (1980): LC-QCD/OPE \Rightarrow



$\pi - \gamma$ Transition Form Factor: $\gamma^*(Q^2) + \pi \rightarrow \gamma$



DSE ladder-rainbow

Triangle Diag

1998 calc

Mitchell, Roberts, PCT



LR: Successes, Problems, Resolutions

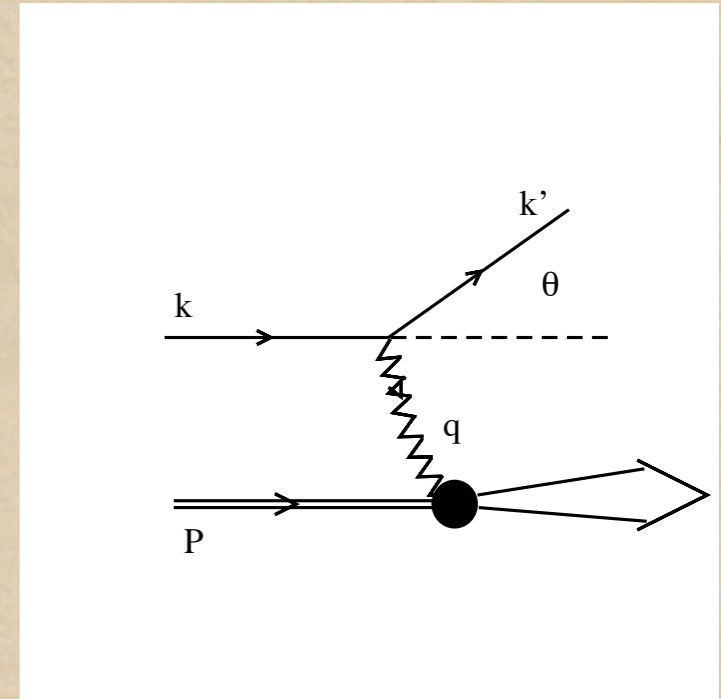
- **Successes:**
 - S-wave mesons, PS and V, light quarks and QQ, no spurious thresholds
 - Exact PS mass formula, Goldstone Thm, ΔM_{HF} from DCSB
 - f_{EW} , strong decays, radiative decays, form factors, $Q^2 < 5\text{GeV}^2$

- **Problems:**
 - Axial vector ($L > 0$) mesons (a_1, b_1, \dots) too light
 - Physical diquarks, no physical V or PS qQ states
 - Excited states are difficult

- **Probable Resolution:**
 - Quark-gluon vertex: $\Gamma_\mu \Rightarrow \Sigma_q \Rightarrow K_{BSE}$
 - Use analysis of spacelike correlators, 3-pt functions

Deep Inelastic Lepton Scattering

- ◆ PDFs: $u_\pi(x)$, $u_K(x)$, $s_\pi(x)$
- ◆ Drell-Yan data exists
- ◆ Pion and Kaon/Pion Ratio
- ◆ Employ LR DSE model
- ◆ Bjorken limit



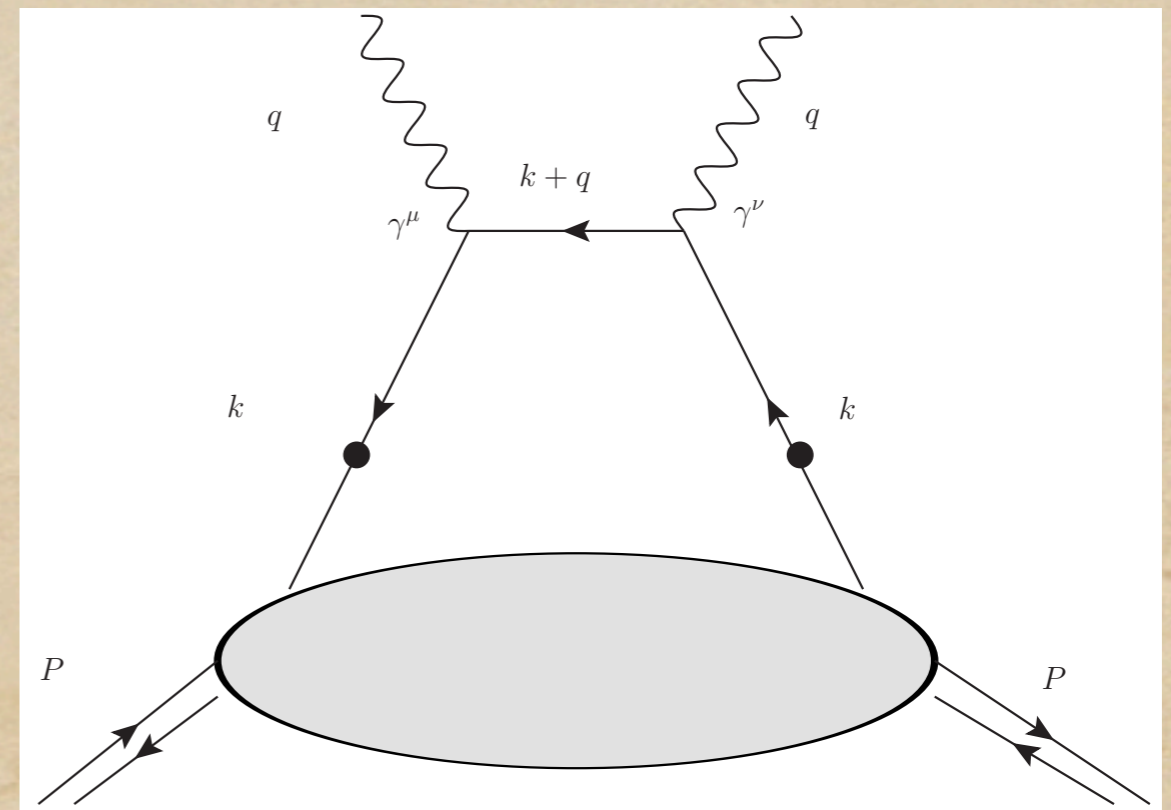
Leading order in OPE

DIS is hard and fast—confinement is soft and slow

$$\Rightarrow S(k+q) \rightarrow \frac{\gamma^+}{2(k^+ - P^+ x) + i\epsilon}$$

$$q^+ = q \cdot n = -P^+ x, \quad |\xi^-| \sim \frac{1}{Mx}$$

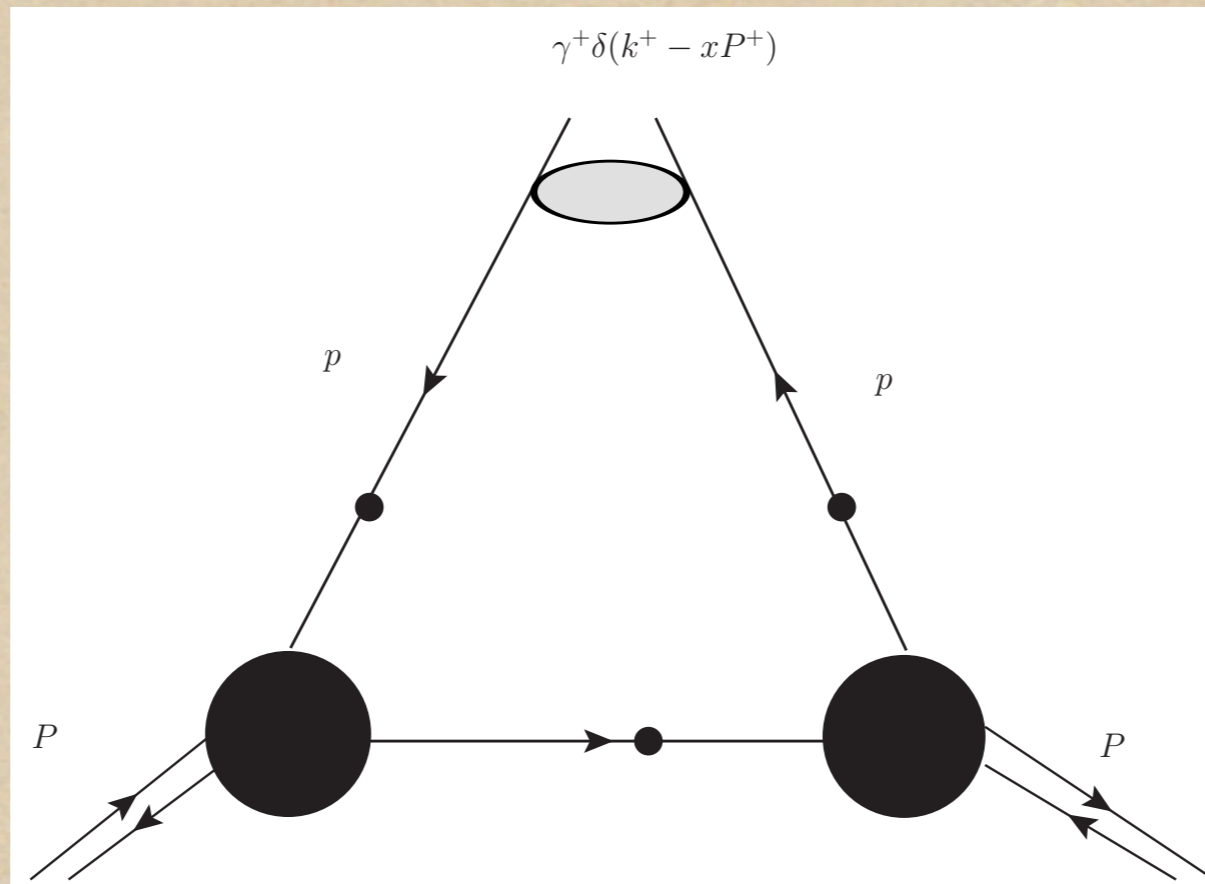
$$q^- = q \cdot p = 2\nu, \quad |\xi^+| \sim 0$$



$$q_f(x) = \frac{1}{4\pi} \int dz^- e^{-ixP^+ z^+} \langle \pi(P) | \bar{\psi}_f(z^-) \gamma^+ \psi_f(0) | \pi(P) \rangle_c = -q_{\bar{f}}(-x)$$

$$N_f^v = \int_0^1 dx [q_f(x) - q_{\bar{f}}(x)] = \frac{1}{2P^+} \langle \pi(P) | J^+(0) | \pi(P) \rangle_c = 1$$

From DSE-BSE at ladder-rainbow truncation



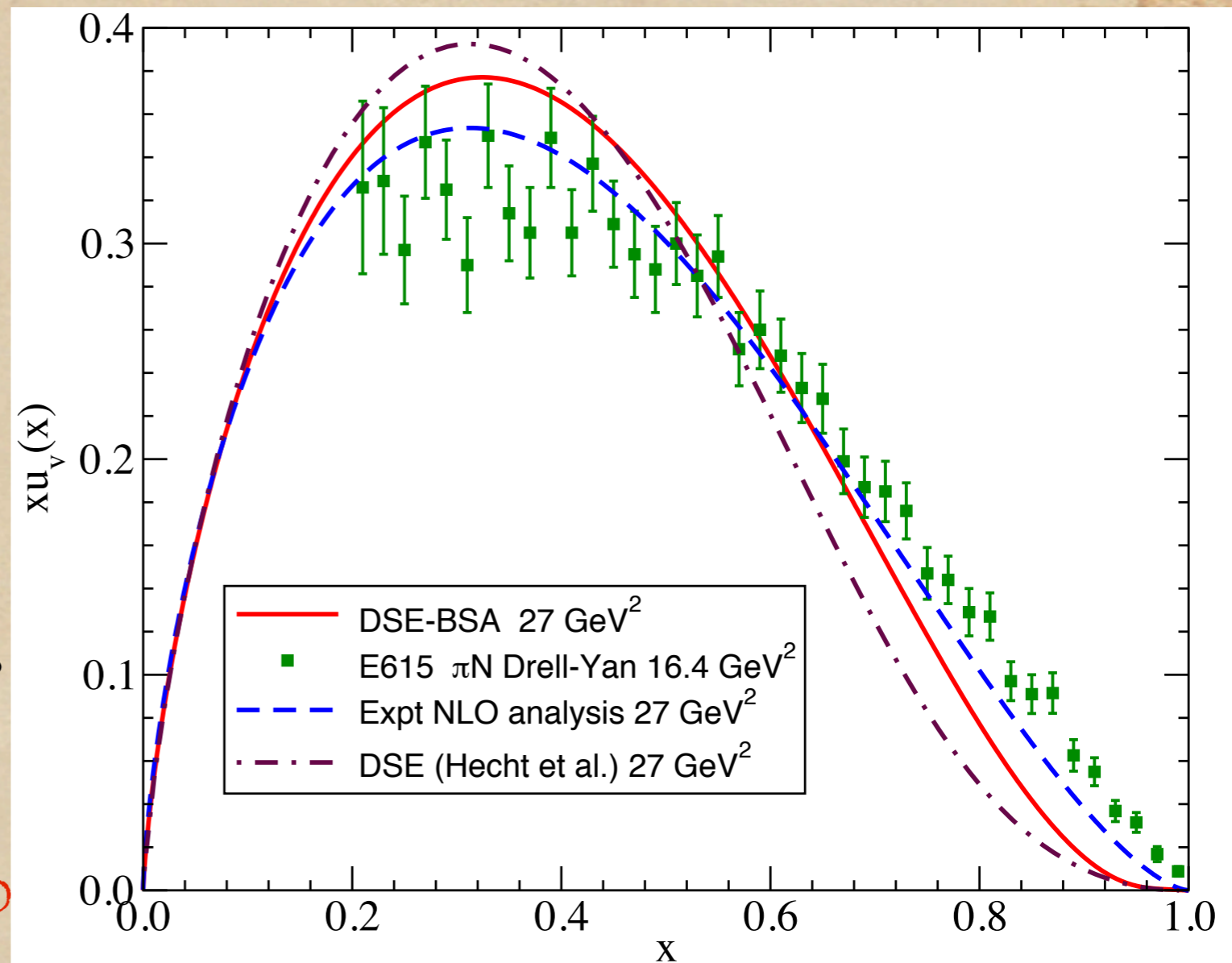
$W^{\mu\nu} \propto \{T^{\mu\nu}(\epsilon) - T^{\mu\nu}(-\epsilon)\} \Rightarrow$ Euclidean model elements can be continued

$$q_f^v(x) = \frac{i}{2} \text{tr}_{\text{cd}} \int_p^\Lambda \Gamma_\pi(p, P) S(p) \Gamma^+(p; x) S(p) \Gamma_\pi(p, P) S(p - P)$$

$$\Gamma^+(p; x) = \gamma^+ \delta(p^+ - xP^+) + \dots$$

Valence $u_\pi(x)$ from DSE-BSE solutions

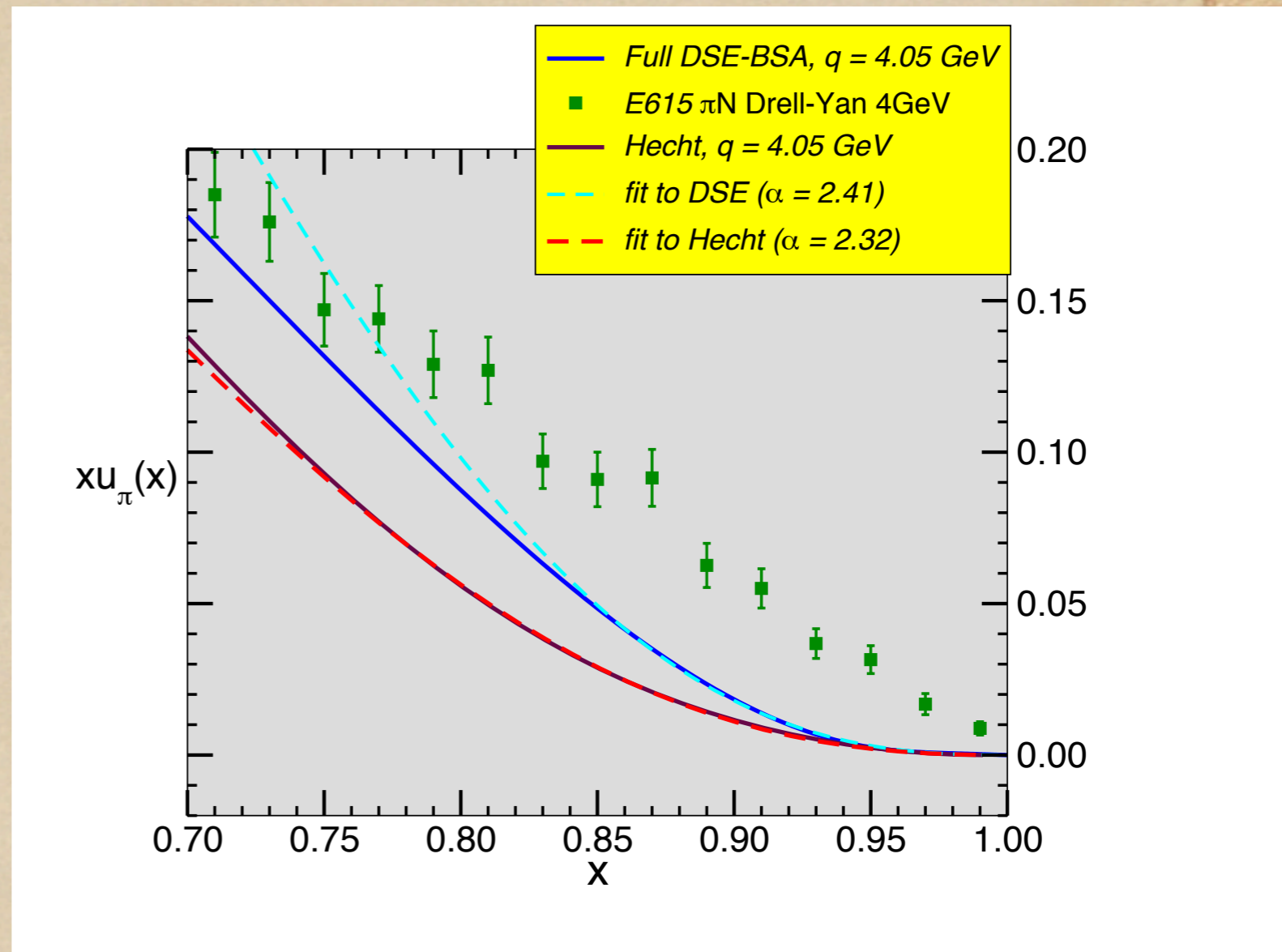
- ◆ Valence quarks, handbag diagram
- ◆ Data: Conway et al, PRD39, 92 (1989). $M_{l\bar{l}} = 4.05$ GeV
- ◆ Prev DSE (phen): Hecht et al, PRC63, 025213 (2001),
 $\Gamma_\pi(k^2, k \cdot P = 0) \sim i\gamma_5 B_0(k^2)/f_\pi^0 + \dots$
 $S_{phen}(k)$
- ◆ Large x behavior: $(1 - x)^\alpha$, $\alpha = ?$
- ◆ T. Nguyen, PhD 2009, KSU,
 Nguyen&PCT, in preparation 2010
- ◆ Wijesooriya, Reimer&Holt, PRC72, 065203 (2005)



Momentum Sum Rule: $\langle x \rangle_{Q_0^2} = 0.76$

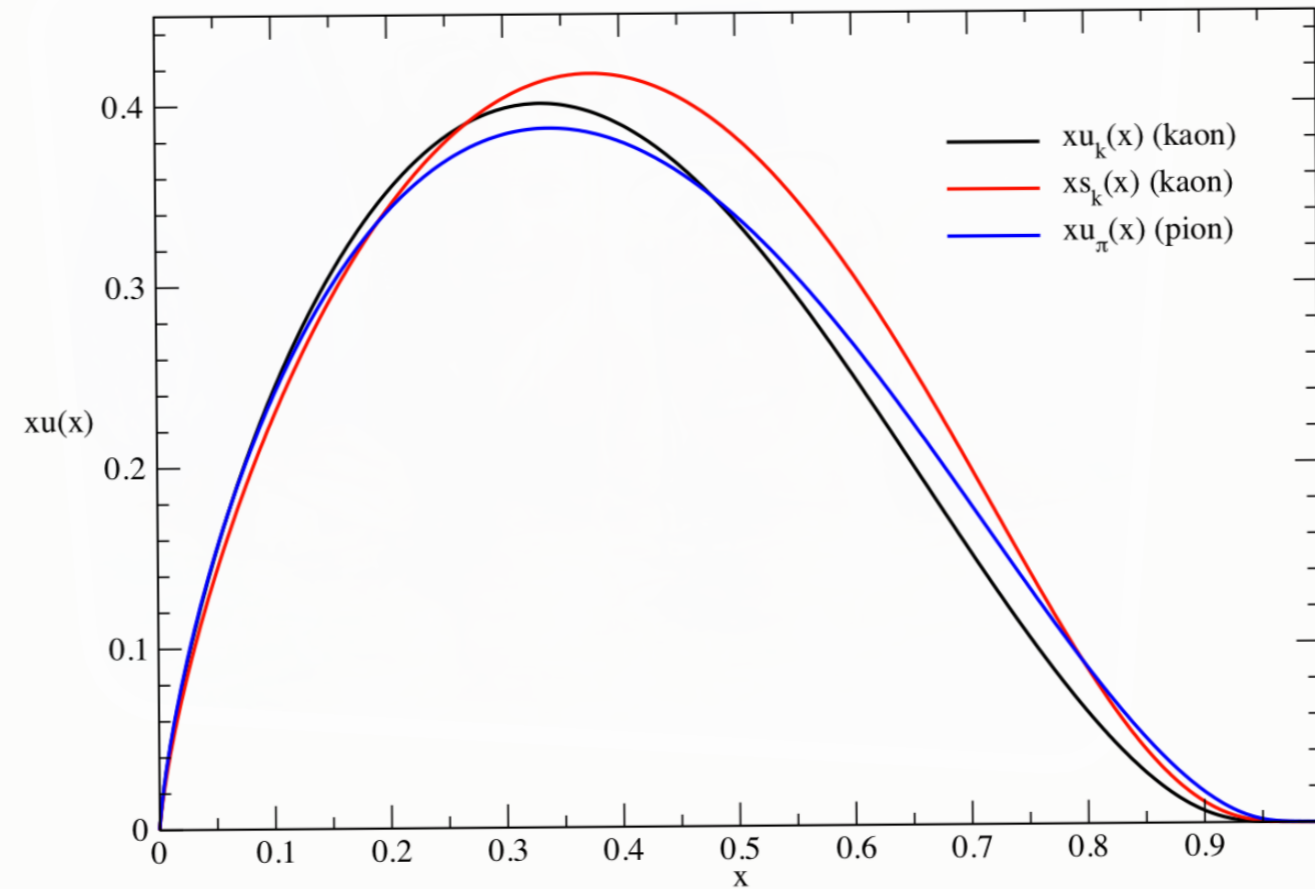
$u_\pi(x)$ at large x ; pQCD

- ◆ Scale for pQCD onset is model-depn.
- ◆ Global DIS fits: $\alpha \sim 1.5$
- ◆ Const. q models, NJL, duality:
 $\alpha \sim 1$
- ◆ pQCD: Farrar-Jackson, Brodsky, Ezawa, DSEs:
 $\alpha = 2 + \gamma(Q^2)$



Quark Distributions in π and K

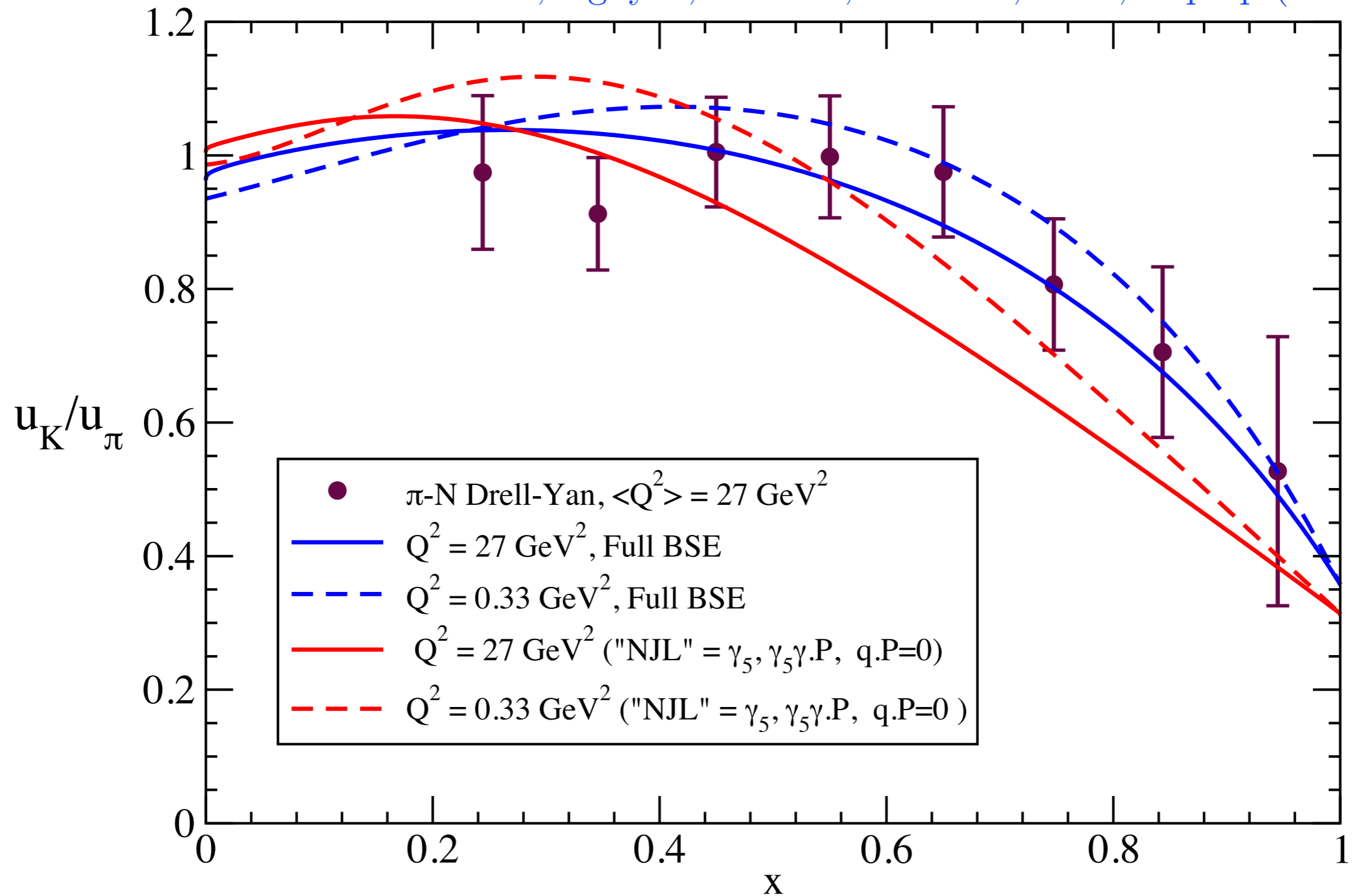
Evolved to $q = 4.05$ GeV



- Environmental depn of $u(x)$ in accordance with effective quark mass
- $u(x)$, $s(x)$ difference in K in accordance with effective quark mass

Environmental Dependence of Valence $u(x)$

—Bashir, Nguyen, Roberts, Souchlas, PCT, in prep (2010)

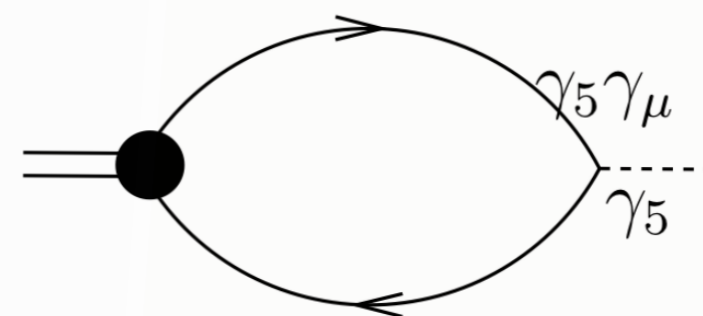
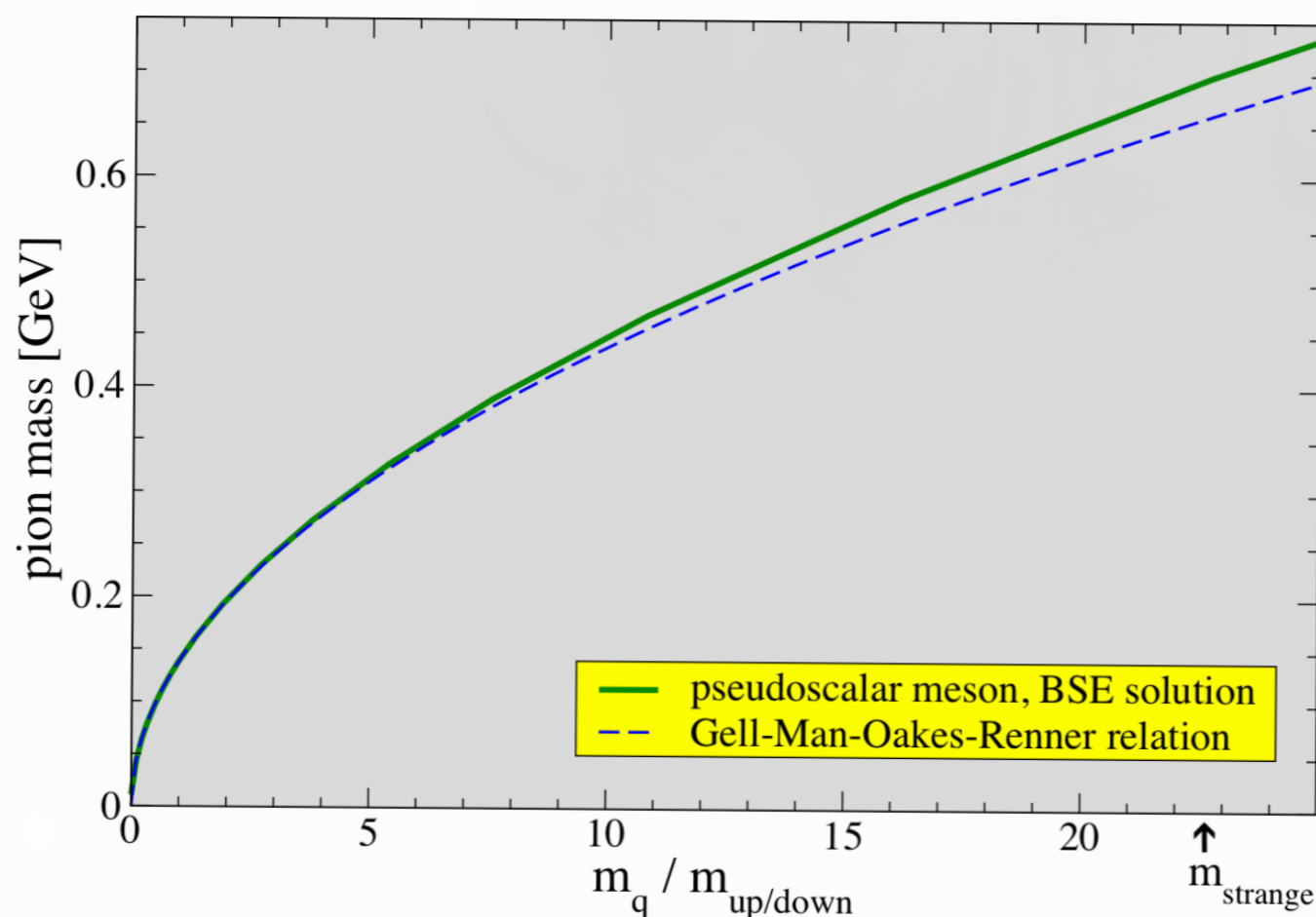


● CERN-SPS data: J. Badier et al, PLB **93**, 354 (1980)



Flavor Non-singlet PS Mass Relation

$$f_H m_H^2 = 2 m_q(\mu) \rho_H(\mu)$$



$$i f_\pi P_\mu = \langle 0 | \bar{q} \gamma_5 \gamma_\mu q | \pi \rangle$$

$$i \rho_\pi = -\langle 0 | \bar{q} i \gamma_5 q | \pi \rangle$$

$$\lim_{m \rightarrow 0} f_\pi \rho_\pi = -\langle \bar{q} q \rangle_\mu$$

$$-\langle \bar{q} q \rangle_\mu^\pi = f_\pi(m) \rho_\pi(m)$$

PM, Roberts, Tandy, PLB420, 267 (1998) [nucl-th/9707003]

In-hadron Condensates

$$- \langle \bar{q}q \rangle_{\mu}^{\pi} = -f_{\pi} \langle 0 | \bar{q} \gamma_5 q | \pi \rangle_{\mu} = f_{\pi}^2 m_{\pi}^2 / 2m(\mu)$$

$$\lim_{m \rightarrow 0} \langle \bar{q}q \rangle_{\mu}^{\pi} = -Z_4(\mu, \Lambda) \text{tr}_{\text{cd}} \int_q^{\Lambda} S_0(q, \mu) = \langle \bar{q}q \rangle_{\mu}^0$$

$\langle \bar{q}q \rangle_{\mu}^0$ is really a property of the PS Goldstone boson BSE wavefunction

Brodsky & Shrock: confinement & DCSB introduce an IR mass scale or max wavelength for virtual fields in hadrons

Brodsky, Roberts, Shrock & PCT, arXiv:1005.4610

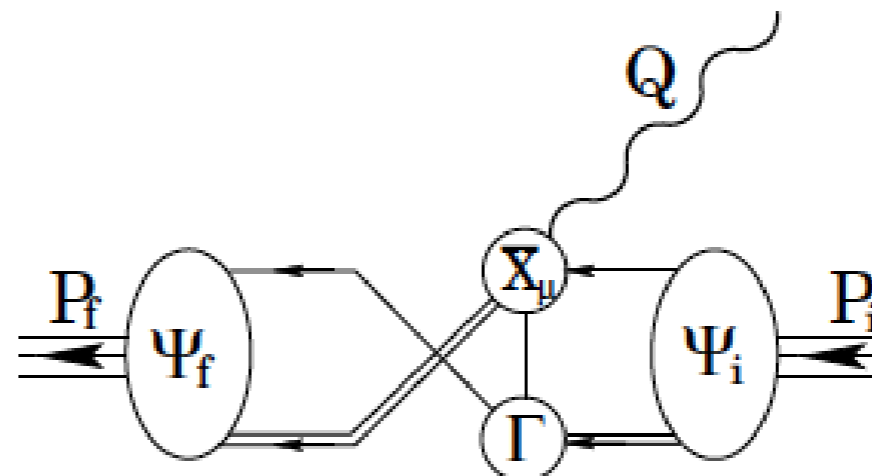
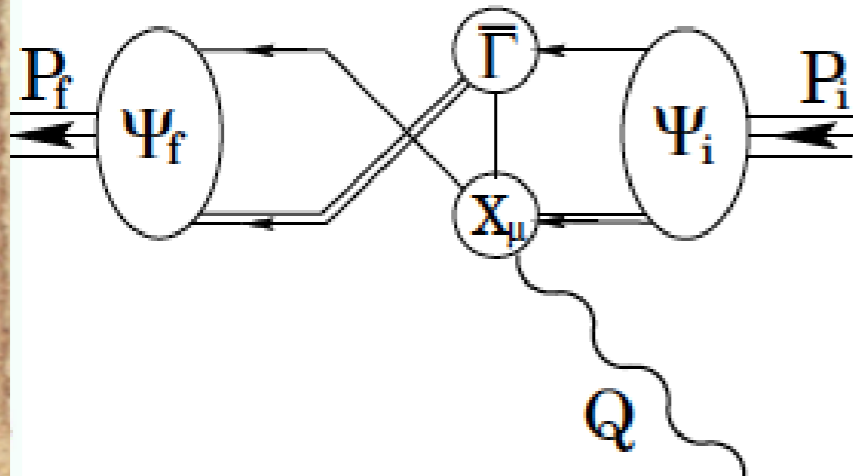
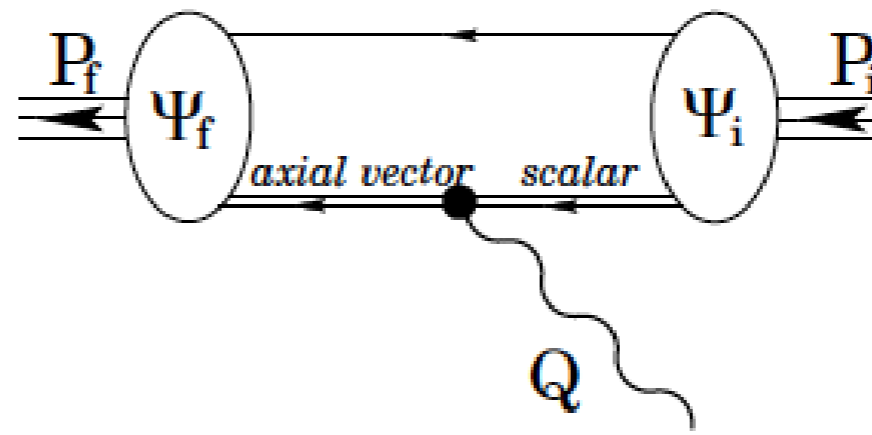
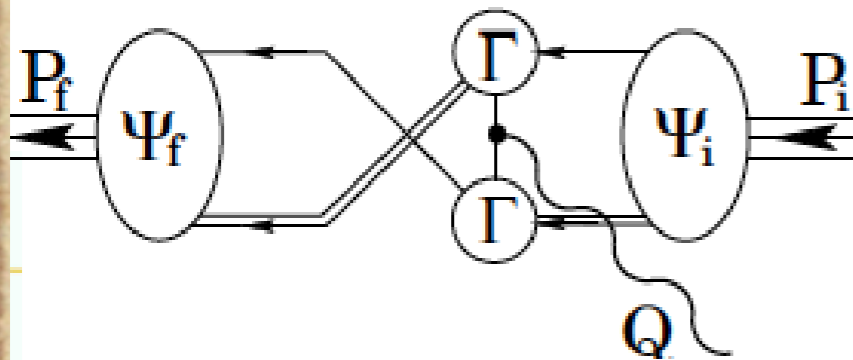
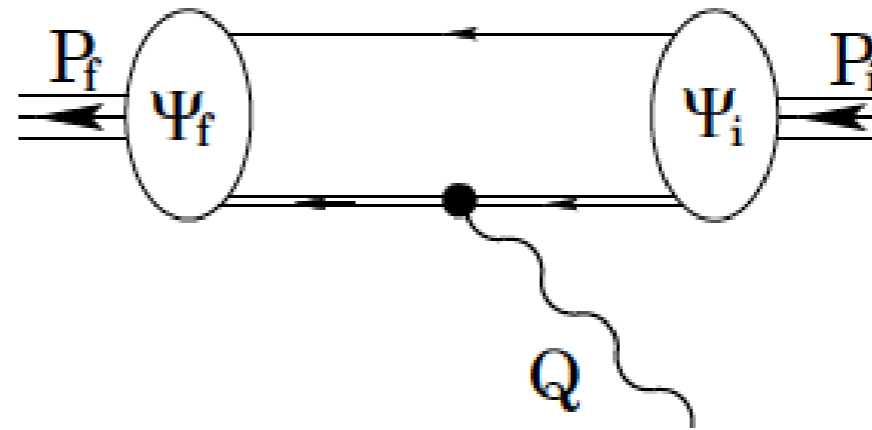
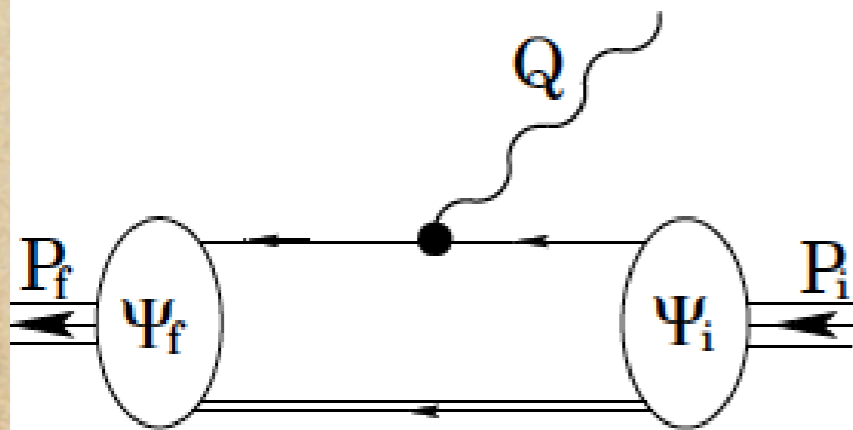
”Essence of the vacuum quark condensate”

Implications for Cosmological Const, and DCSB in Light-Front Field Theory

Nucleon-Photon Vertex

6 terms ...

constructed systematically ... current conserved automatically
 for on-shell nucleons described by Faddeev Amplitude





Cloët, Roberts *et al.*

– arXiv:0710.2059 [nucl-th]

– arXiv:0710.5746 [nucl-th]

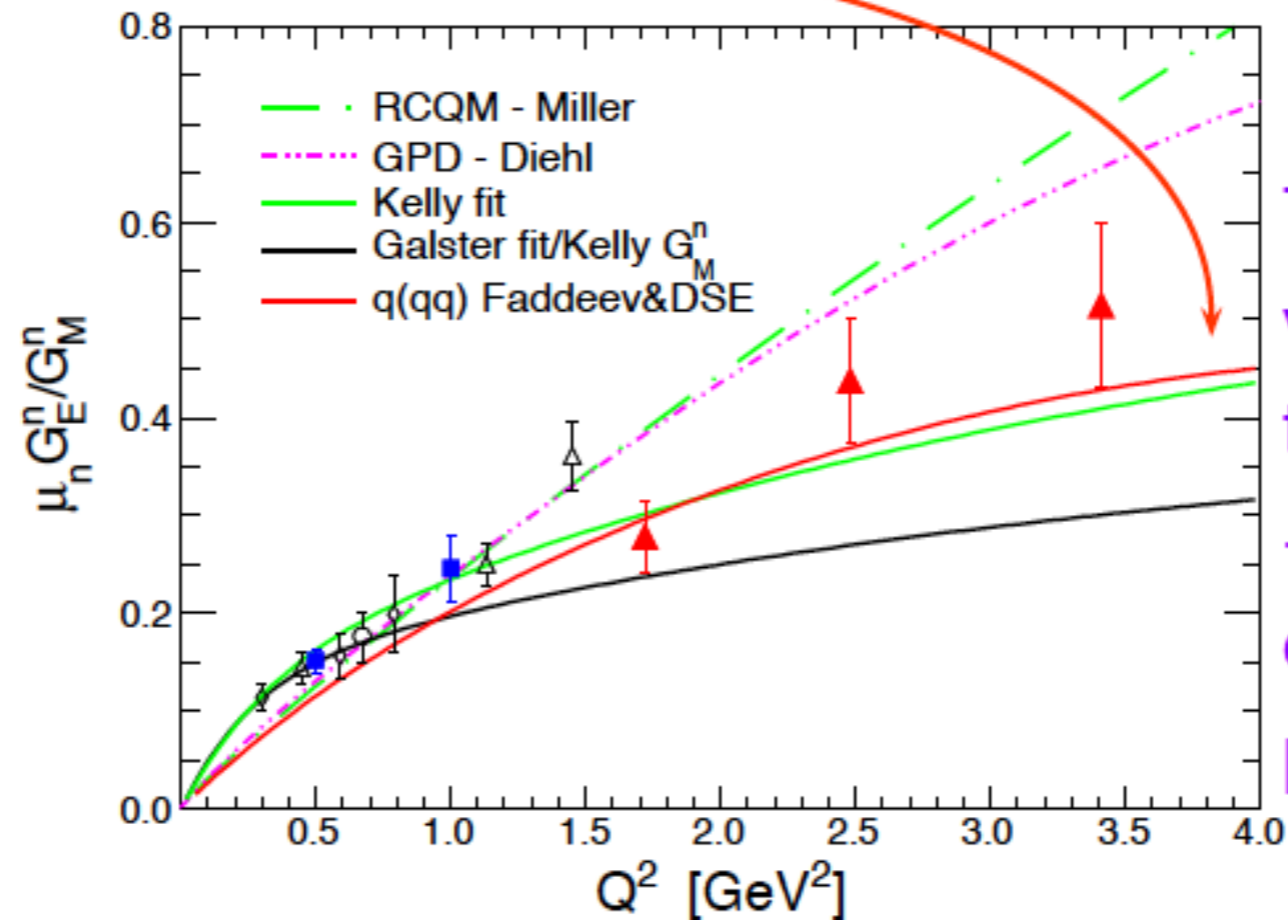
– arXiv:0804.3118 [nucl-th]

– arXiv:0812.0416 [nucl-th] – *Survey of nucleon EM form factors*

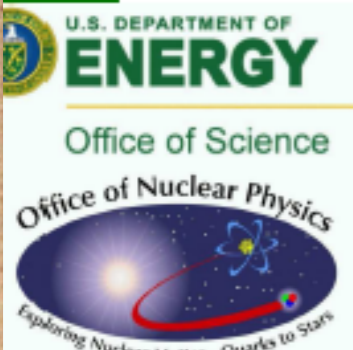
$$\frac{\mu_n G_E(Q^2)}{G_M(Q^2)}$$

● DSE-Faddeev Equation prediction

Red solid curve



This evolution very sensitive to momentum-dependence dressed-quark propagator



B. Wojtsekhowski, Jefferson Lab E02-013 Collaboration, *in preparation.*

Figure courtesy S. Riordan



Cloët, Roberts *et al.*

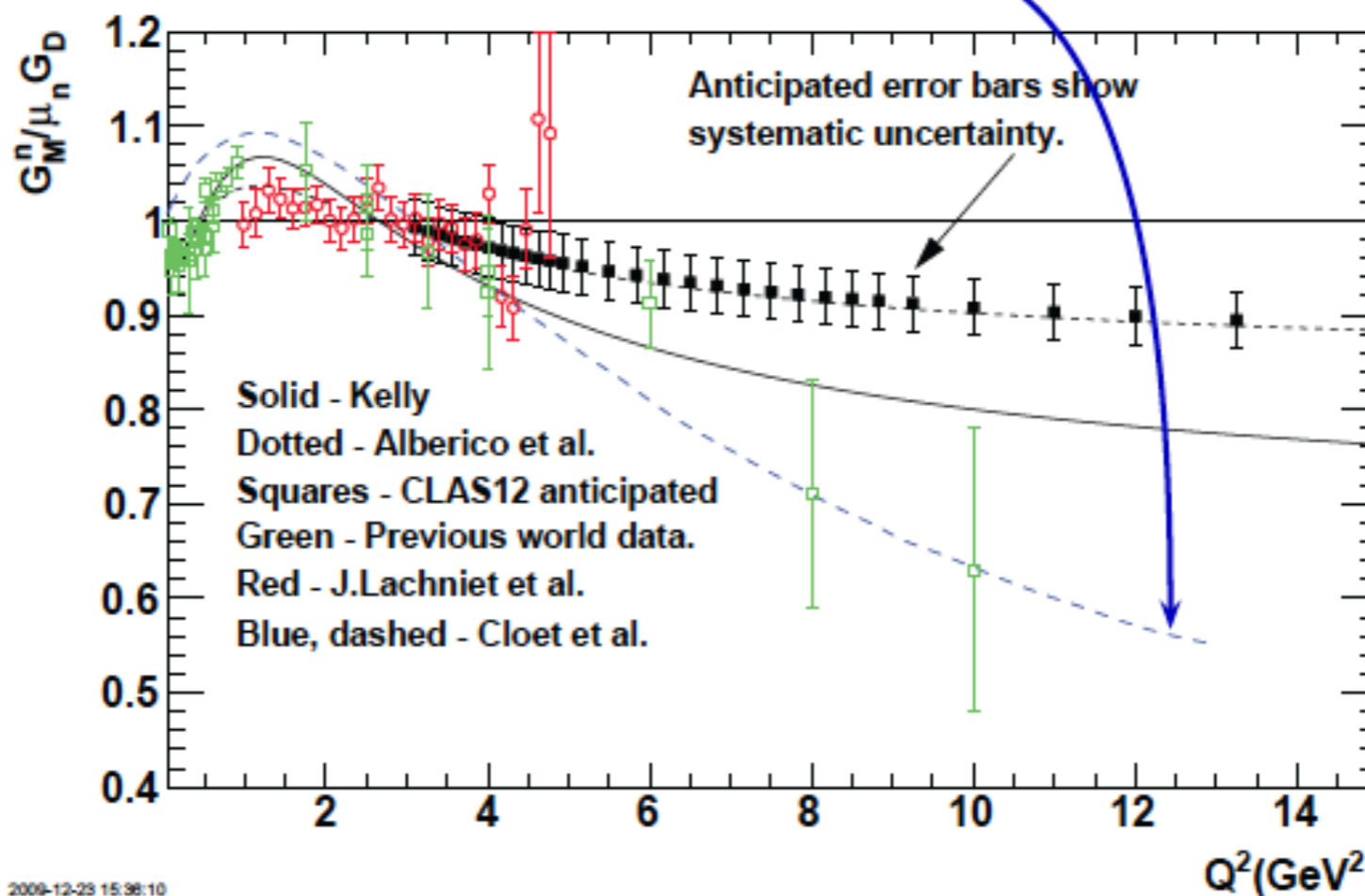
- arXiv:0710.2059 [nucl-th]
- arXiv:0710.5746 [nucl-th]
- arXiv:0804.3118 [nucl-th]

- arXiv:0812.0416 [nucl-th] - *Survey of nucleon EM form factors*

$$\frac{G_M^n(Q^2)}{\mu_n G_D(Q^2)}$$

● DSE-Faddeev Equation prediction

Blue long-dashed curve

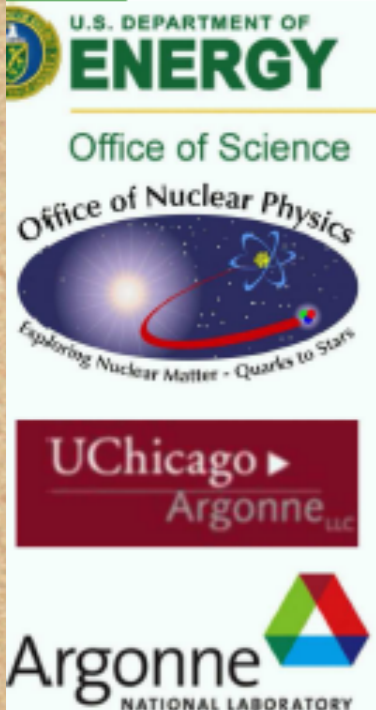


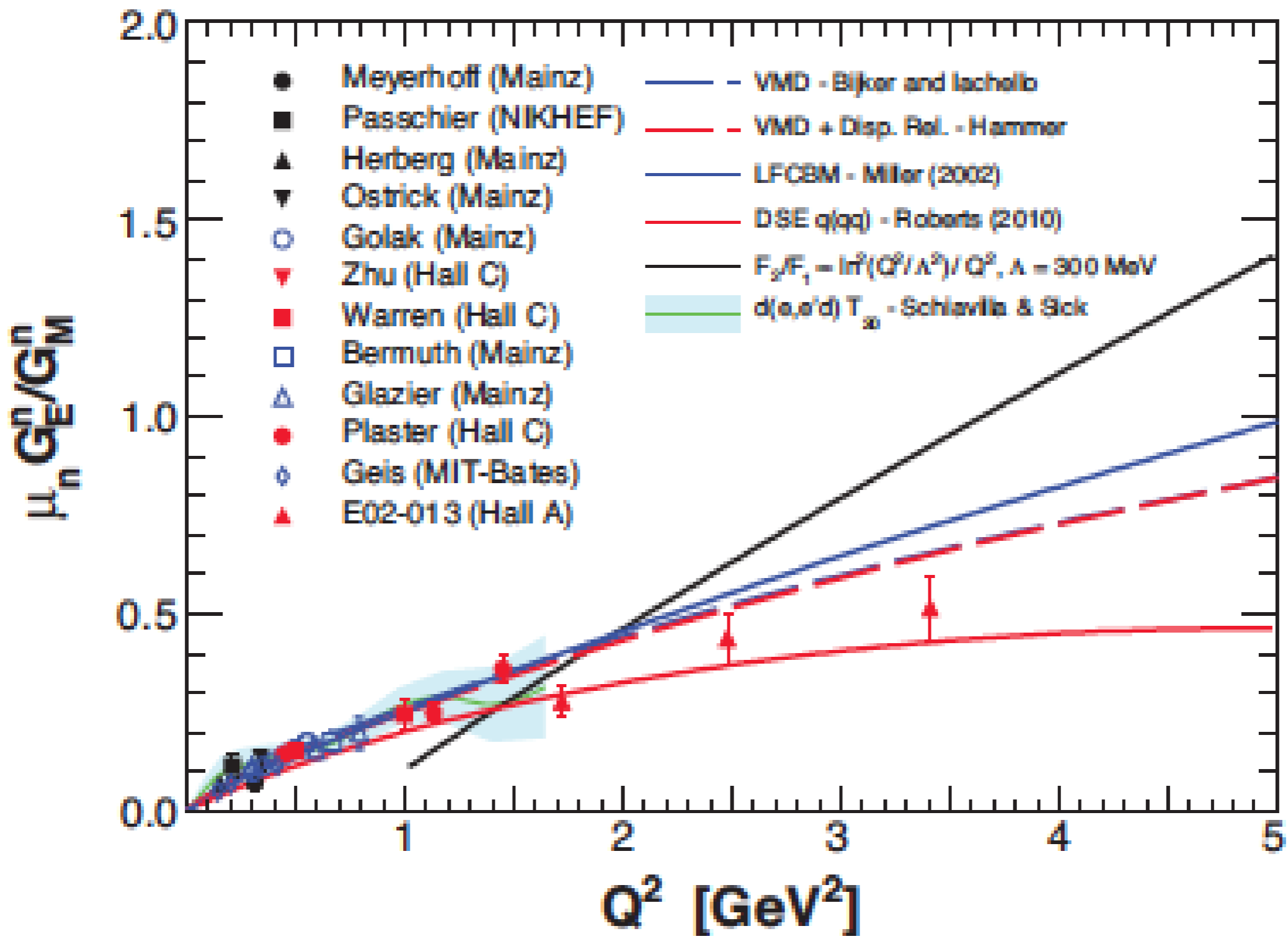
Sensitivity to $M(p^2)$ means experiments probe IR behaviour of strong running coupling

2009-12-29 15:38:10

Jefferson Lab E12-07-104, 12GeV Proposal.

Gilfoyle, Brooks, Hafidi for CLAS Collaboration





JLab, May 19, 2010, 20



Axial anomaly and $\eta - \eta'$ states

- Ch symm: $\partial_\mu(z) \langle j_{5\mu}^\alpha(z) q(x) \bar{q}(y) \rangle$ involves $2 \text{tr}_f(\mathcal{F}^\alpha) \langle Q_t(z) q(x) \bar{q}(y) \rangle$
- Matrix elements, amputated \Rightarrow AV-WTI

$$P_\mu \Gamma_{5\mu}^\alpha(k; P) = -2i \mathcal{M}^{\alpha\beta} \Gamma_5^\beta(k; P) - \delta_{\alpha,0} \Gamma_A(k; P) + S^{-1}(k_+) i\gamma_5 \mathcal{F}^\alpha + i\gamma_5 \mathcal{F}^\alpha S^{-1}(k_-)$$

- Residues at PS poles \Rightarrow PS mass formula for arbitrary m_q , any flavor:

$$m_p^2 f_p^\alpha = 2 \mathcal{M}^{\alpha\beta} \rho_p^\beta + \delta^{\alpha,0} n_p, \quad n_p = 2 \text{tr}_f(\mathcal{F}^0) \langle 0|Q_t|p\rangle$$

$$\rho_p^\alpha(\mu) = \langle 0|\bar{q} \gamma_5 \mathcal{F}^\alpha q|p\rangle, \quad p = \text{any PS}$$

—[Bhagwat, Chang, Liu, Roberts, PCT, PRC (76), 2007; arXiv:0708.1118]

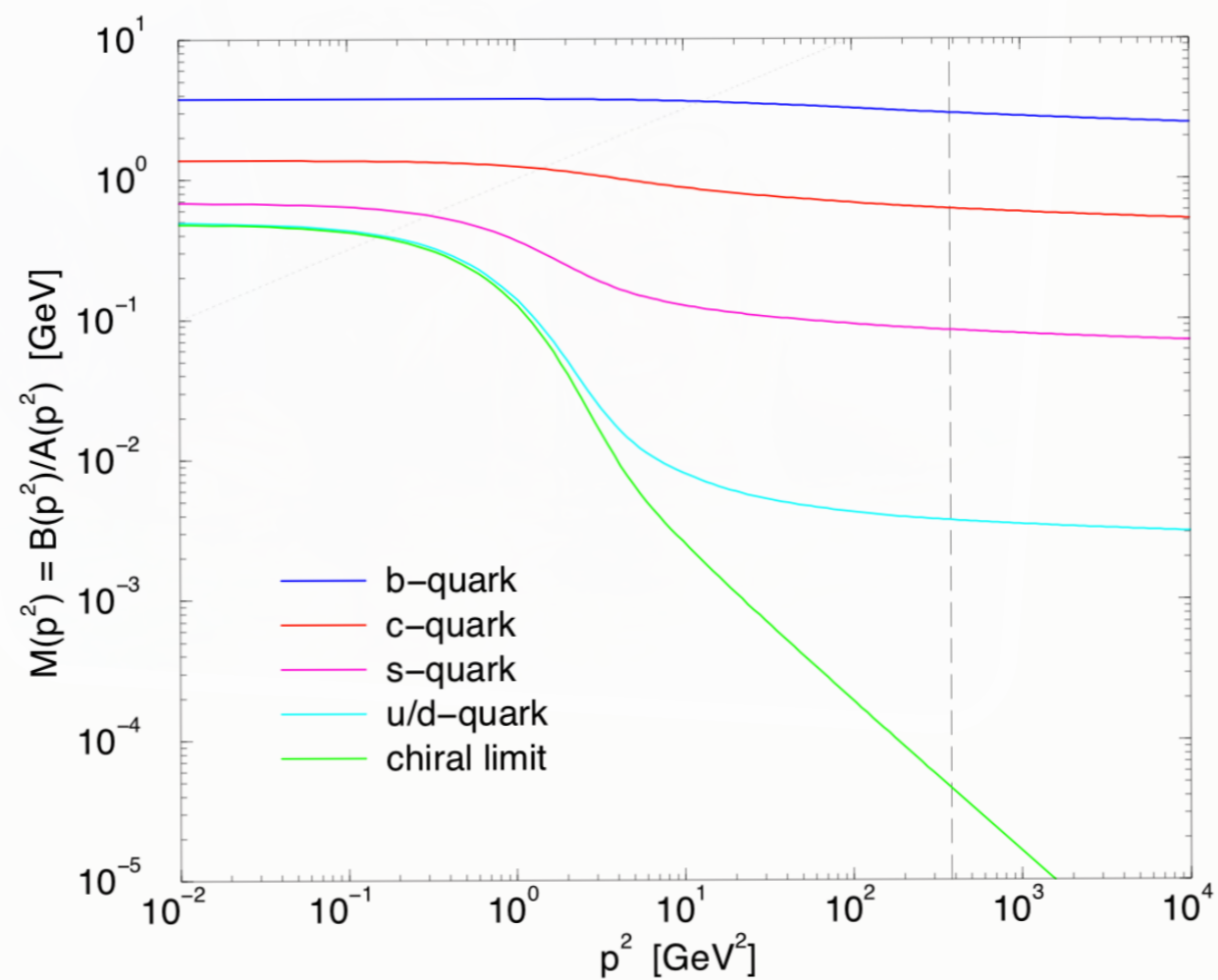


$\pi^0 - \eta - \eta'$ mixing: 3 flavors

- $m_u - m_d$ causes π^0 to be mixed in:
135 MeV : $|\pi^0\rangle \sim 0.72 \bar{u}u - 0.69 \bar{d}d - 0.013 \bar{s}s$
455 MeV : $|\eta\rangle \sim 0.53 \bar{u}u + 0.57 \bar{d}d - 0.63 \bar{s}s$
922 MeV : $|\eta'\rangle \sim 0.44 \bar{u}u + 0.45 \bar{d}d + 0.78 \bar{s}s$
- $m_u = m_d \Rightarrow$
455 MeV : $|\eta\rangle \sim 0.55 (\bar{u}u + \bar{d}d) - 0.63 \bar{s}s, \quad \theta_\eta = -15.4^\circ$
924 MeV : $|\eta'\rangle \sim 0.45 (\bar{u}u + \bar{d}d) + 0.78 \bar{s}s, \quad \theta_{\eta'} = -15.7^\circ$
- Chiral limit: $m_{\eta'}^2 = (0.852 \text{ GeV})^2 \equiv 2\text{tr}_f(\mathcal{F}^0) \langle 0|Q_t|\eta'\rangle / f_{\eta'}^0$
- cf Witten-Veneziano a-v ghost scenario $\Rightarrow m_{\eta'}^2 = h^2 + m_{\text{GB}}^2$
- It is worth extending to a realistic LR model for K_N with separable K_A : one obtains access to decay constants, residues, and details of the mass relations



Quark mass functions from DSE solutions





Constituent Mass Concept for *c*- and *b*-quarks

All GeV	D(uc)	D*(uc)	D _s (sc)	D* _s (sc)
expt M	1.86	2.01	1.97	2.11
calc M	1.85(FIT)	2.04	1.97	2.17
expt f	0.222	?	0.294	?
calc f	0.154	0.160	0.197	0.180

All GeV	B(ub)	B*(ub)	B _s (sb)	B* _s (sb)	B _c (cb)	B* _c (cb)
expt M	5.28	5.33	5.37	5.41	6.29	?
calc M	5.27(FIT)	5.32	5.38	5.42	6.36	6.44
expt f	0.176	?	?	?	?	?
calc f	0.105	0.182	0.144	0.20	0.210	0.18

- Fit \Rightarrow constituent masses: $M_c^{\text{cons}} = 2.0 \text{ GeV}$, $M_b^{\text{cons}} = 5.3 \text{ GeV}$
- Consistent with $M^{DSE}(p^2 \sim -M^2)$ generated by $m_c = 1.2 \pm 0.2$, $m_b = 4.2 \pm 0.2$, [PDG, $\mu = 2 \text{ GeV}$]
- Does heavy quark dressing contribute anything? Too much in this DSE model—no mass shell !



Quarkonia

All GeV	M_{η_c}	f_{η_c}	$M_{J/\psi}$	$f_{J/\psi}$
expt	2.98	0.340	3.09	0.411
calc with M_c^{cons}	3.02	0.239	3.19	0.198
calc with $\Sigma_c^{\text{DSE}}(p^2)$	3.04	0.387	3.24	0.415

All GeV	M_{η_b}	f_{η_b}	M_{Υ}	f_{Υ}
expt	9.4 ?	?	9.46	0.708
calc with M_b^{cons}	9.6	0.244	9.65	0.210
calc with $\Sigma_b^{\text{DSE}}(p^2)$	9.59	0.692	9.66	0.682

- QQ and qQ decay constants too low by 30-50% in **constituent mass approximation**
- Quarkonia decay constants much better for **DSE** dressed quarks (within 5% of expt.)
- IR sector (gluon k below ~ 0.8 GeV) contribute little for bb or cc quarkonia in DSE, BSEs
- QQ states are more point-like than qq or qQ states



Recovery of a qQ Mass Shell

- Suppress gluon k below ~ 0.8 GeV in DSE dressing of b propagator
- Retain IR sector for dressed "light" quark and BSE kernel
- Now a mass shell is produced

All GeV	B(ub)	B*(ub)	B _s (sb)	B* _s (sb)	B _c (cb)	B* _c (cb)
expt M	5.28	5.33	5.37	5.41	6.29	?
calc M	4.66	—	4.75	—	5.83	—
expt f	0.176	?	?	?	?	?
calc f	0.133	—	0.164	—	0.453	—

- Masses are ~ 10 % low
- It makes sense that $R_b < R_{qQ} \Rightarrow$ greater limit on low k in Σ_b
- May be partial confirmation of Brodsky and Shrock's suggestion of universal maximum wavelength for quarks/gluons in hadrons [Phys. Lett. B666, (2008)]

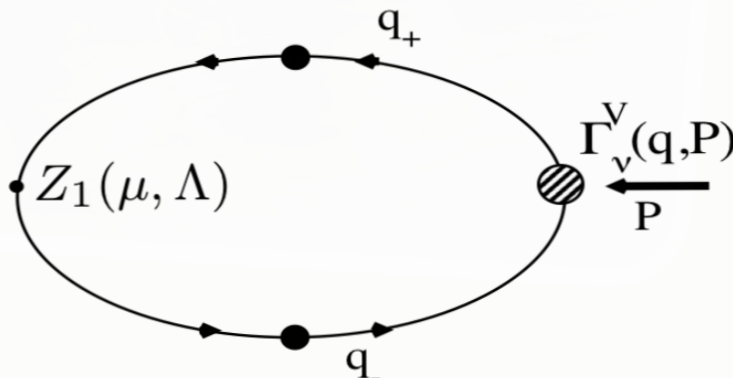


The V-A Current Correlator

- $\Pi_{\mu\nu}^V(x) = \langle 0 | T j_\mu(x) j_\nu^\dagger(0) | 0 \rangle$, isovector currents $j_\mu = \bar{u}\gamma_\mu d$, $j_\mu^5 = \bar{u}\gamma_5\gamma_\mu d$

$$\Pi_{\mu\nu}^V(P) = (P^2\delta_{\mu\nu} - P_\mu P_\nu) \Pi^V(P^2)$$

$$\Pi_{\mu\nu}^A(P) = (P^2\delta_{\mu\nu} - P_\mu P_\nu) \Pi^A(P^2) + P_\mu P_\nu \Pi^L(P^2)$$

$$\Pi_{\mu\nu}^V(P) = - \int_q^\Lambda \gamma_\mu \cdot Z_1(\mu, \Lambda) \cdot \Gamma_\nu^V(q, P)$$


- $m_q = 0$: $\Pi^V - \Pi^A = 0$, to all orders in pQCD
- $\Pi^V - \Pi^A$ probes the scale for onset of non-perturbative phenomena in QCD



Physics from the V-A correlator:

OPE:

$$\Pi^{V-A}(P^2) = \frac{32\pi\alpha_s \langle \bar{q}q\bar{q}q \rangle}{9 P^6} \left\{ 1 + \frac{\alpha_s}{4\pi} \left[\frac{247}{4\pi} + \ln\left(\frac{\mu^2}{P^2}\right) \right] \right\} + \mathcal{O}\left(\frac{1}{P^8}\right)$$

Model	$-\langle \bar{q}q \rangle_{\mu=19} (GeV)^3$	$\langle \bar{q}q\bar{q}q \rangle_{\mu=19} (GeV)^6$	$R(\mu=19)$
LR DSE	$(0.216)^3$	$(0.235)^6$	1.65

Weinberg et al Sum Rules:

- I: $\frac{1}{4\pi^2} \int_0^\infty ds [\rho_v(s) - \rho_a(s)] = [P^2 \Pi^{V-A}(P^2)]_{P^2 \rightarrow 0} = -f_\pi^2$
- II: $P^2 [P^2 \Pi^{V-A}(P^2)]|_{P^2 \rightarrow \infty} = 0$
- DGMLY: $\int_0^\infty dP^2 [P^2 \Pi^{V-A}(P^2)] = -\frac{4\pi f_\pi^2}{3\alpha} [m_{\pi^\pm}^2 - m_{\pi^0}^2]$

Model	$f_\pi^2 (GeV^2)$	$f_\pi (MeV)$	f_π^{exp}/f_π^{num}	$\Delta m_\pi (MeV)$	$(\Delta m_\pi)_{exp}$
LR DSE	0.0081	90.0	1.03	4.88	4.43 ± 0.03





Summary

- Effective ladder-rainbow model based on QCD -DSEs; $\langle \bar{q}q \rangle_\mu \Rightarrow 1$ IR parameter
- Convenient and covariant approach to hadronic form factors: N, π , various transitions
- Ground state qQ and QQ mesons (V & PS) up to b-quark region
- Dynamical dressing in $S(p)$ at each stage increases the value of the decay constant [factor of 3 for $\bar{b}b$, factor of 2 for $\bar{c}c$] !
- First combination of BSE-DSE solutions for pion and kaon DIS distributions $u(x), s(x)$
- Used $\langle J J \rangle$, V-A, to estimate $\langle \bar{q}q\bar{q}q \rangle$ as $\sim 70\%$ greater than vac saturation, and npQCD enters at scale 0.5 fm.



Collaborators

- Craig Roberts, Argonne National Lab
- Pieter Maris, Iowa State University
- Yu-xin Liu, Lei Chang, Peking University
- Nick Souchlas, Trang Nguyen, Kent State University

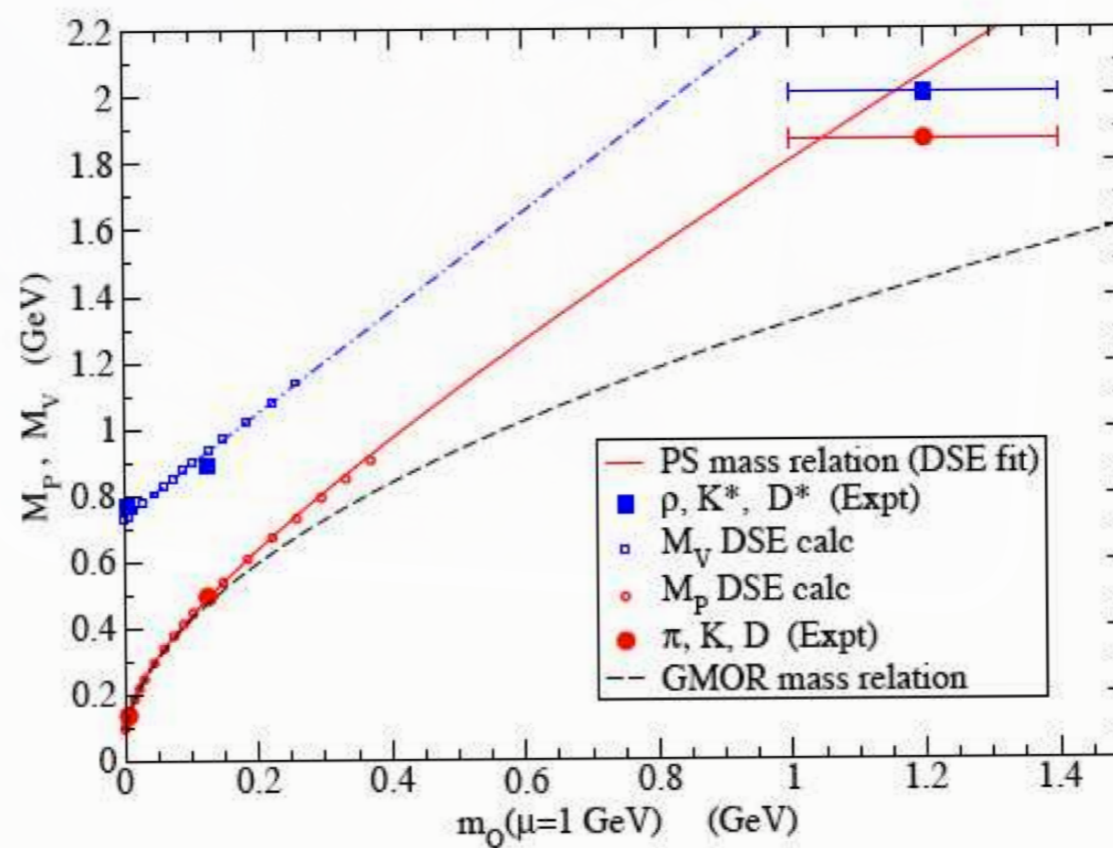
Thankyou!





Inaccuracy of GMOR

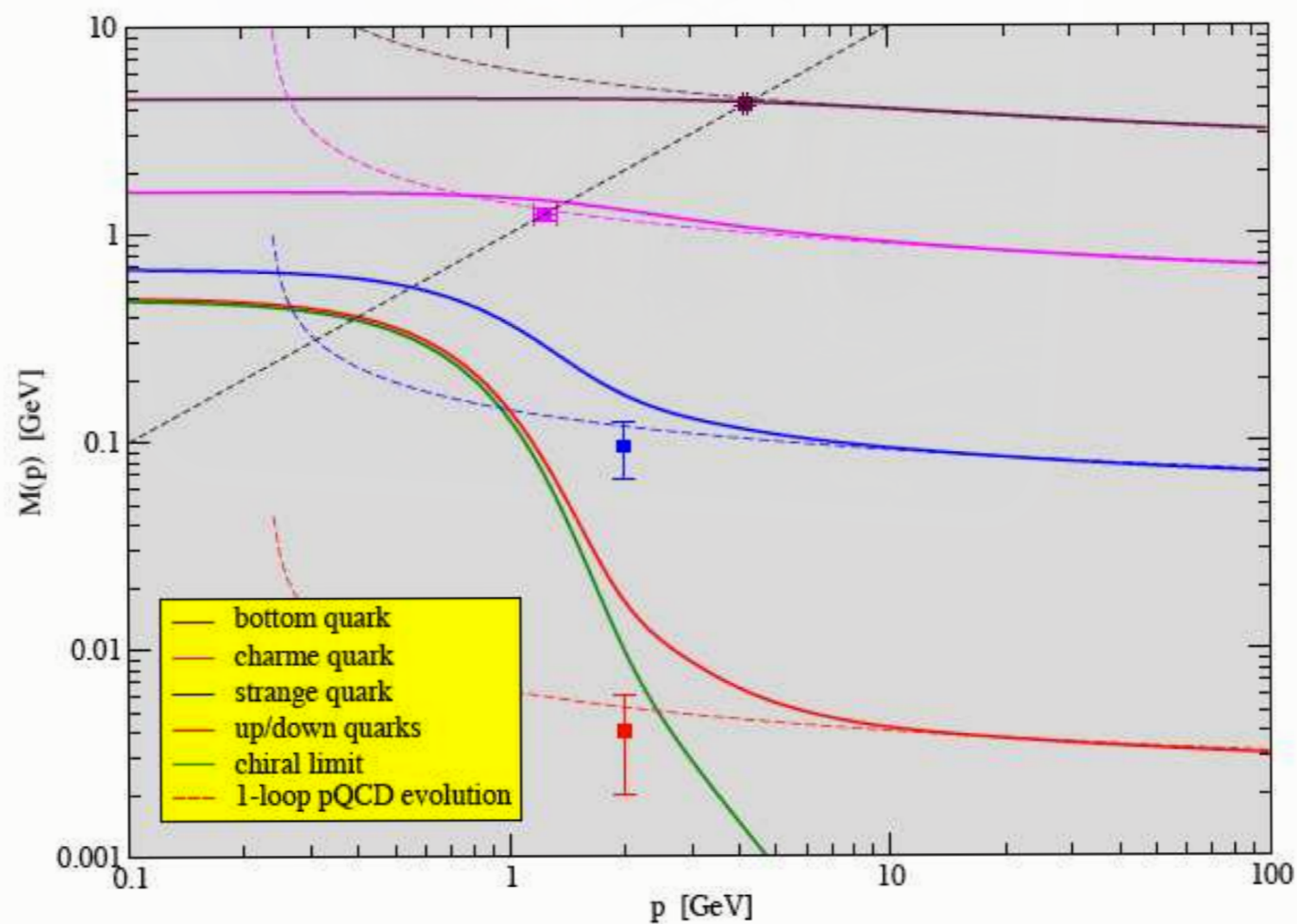
qQ case:



GMOR: 0.2%(π); 4%(K); 14%(0.4GeV); 30%(D)



Compare Quark Masses with PDG





From Gluon vertex to BSE Kernel

- A symmetry-preserving procedure [Bender, Roberts, von Smekal, PLB380, (1996), nucl-th/9602012; Munczek 1995] ; Axial vector and vector WTIs, and Goldstone Thm preserved

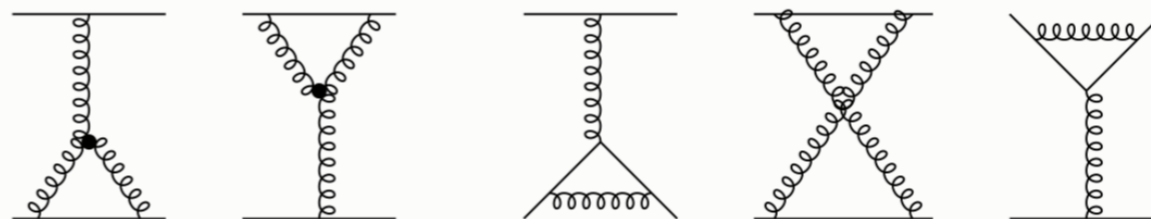
- $K_{\text{BSE}}(x', y'; x, y) = -\frac{\delta}{\delta S(x, y)} \Sigma(x', y')$

- Vertex $\Gamma_\mu(p, q) = \sum \text{diagrams} \Rightarrow K_{\text{BSE}} = \sum \text{diagrams}$

- If Σ contains:



- K_{BSE} contains:

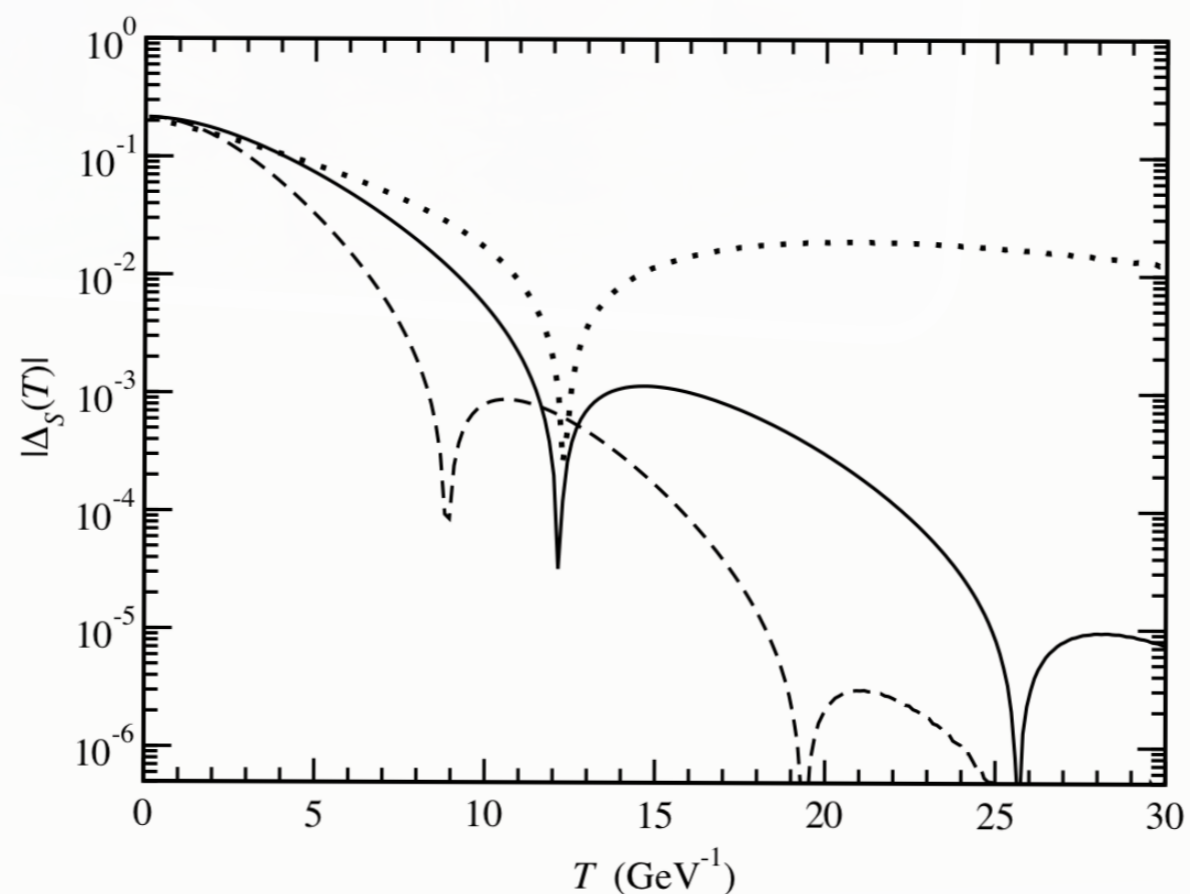


- Independent of model parameters. Model does not fight chiral symmetry, use light vector mesons to fix parameters



Quark Confinement—positivity violation

- Confinement/positivity analysis (Osterwalder-Schrader axiom No. 3)
- Fourier transf $\sigma_S(p_4, \vec{p} = 0)$ to Eucl time T



solid = lattice prop, dashed = MT DSE, dotted = cc pole eq



DSE kernel constrained from Lattice QCD

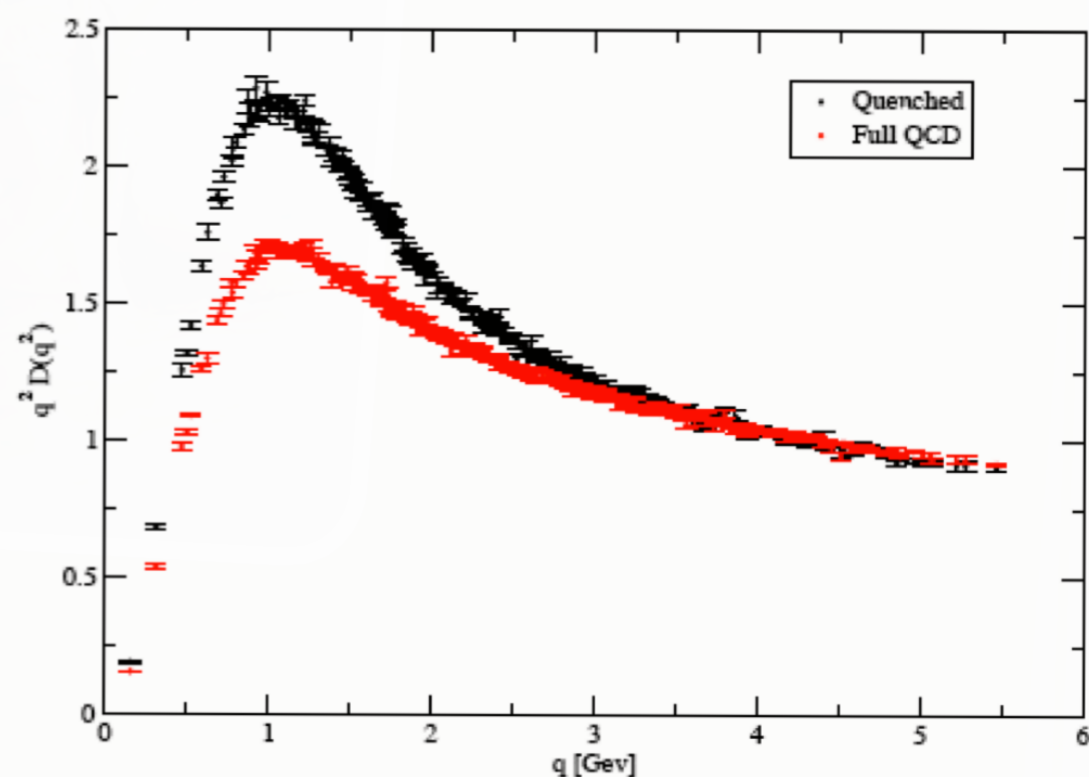
— Bhagwat, Pichowsky, Roberts, Tandy, PRC68, 015203 (03)

- Qu-lattice $D_{\text{gluon}}(q)$

Leinweber, Bowman et al
PRD60, hep-lat/9811027

- Find $\Gamma_{\nu}^{\text{eff}}(q, p)$ so DSE produces

$S_{\text{latt}}(p)$ from $D_{\text{latt}}(q)$



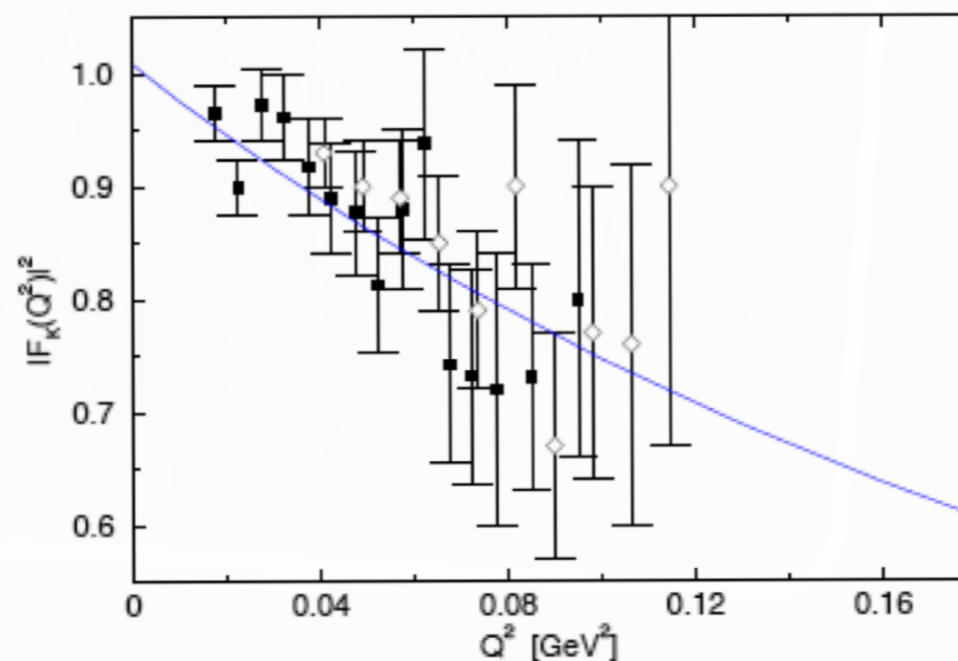
$$g^2 \gamma_{\mu} D(p - q) Z_{1F}(\mu, \Lambda) \Gamma_{\nu}(q, p) \rightarrow \gamma_{\mu} g^2 D(p - q) \gamma_{\nu} V(p - q)$$

UV limit: $g^2 D(k^2) V(k^2) \rightarrow \frac{4\pi\alpha_s^{1-\text{loop}}(k^2)}{k^2}$



Kaon $F(Q^2)$: Low Q^2

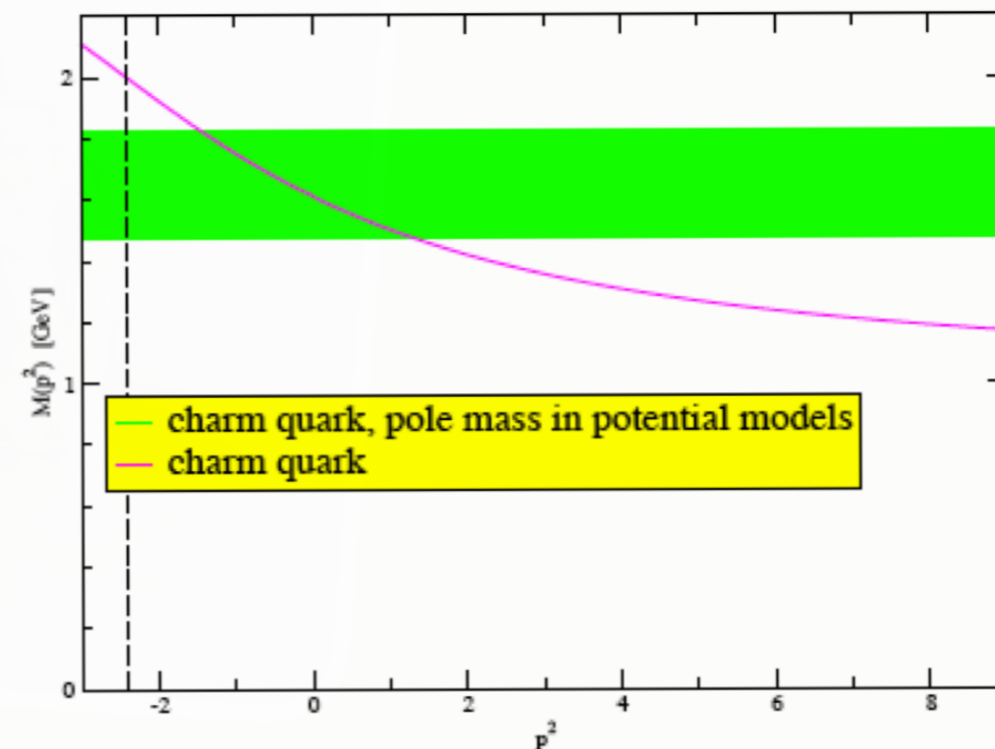
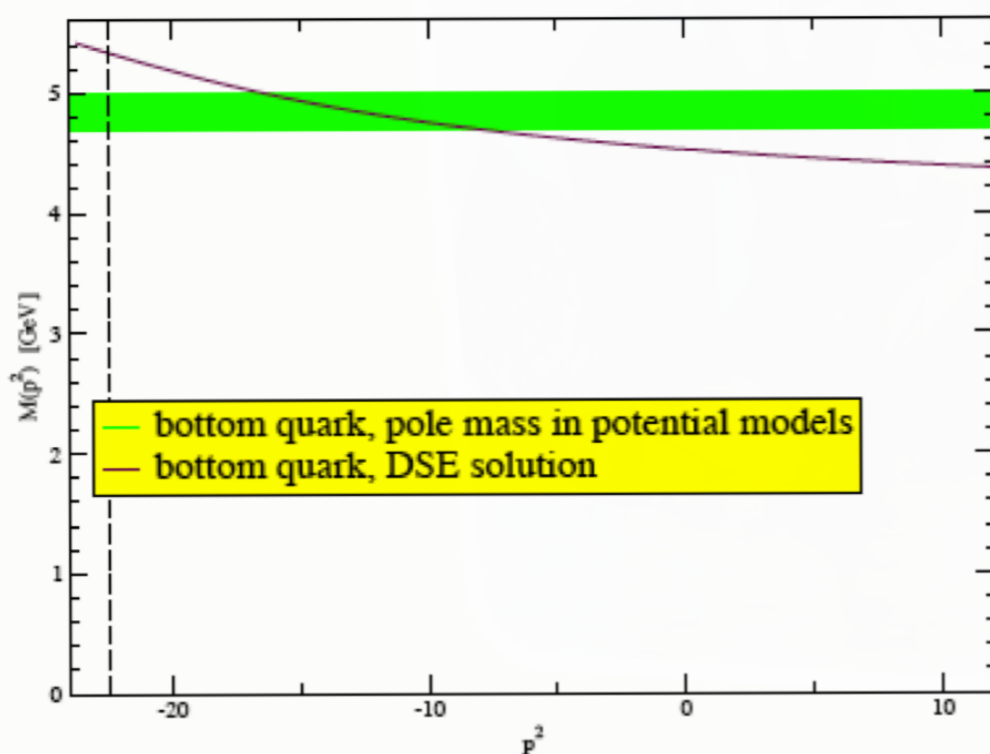
- Impulse approx + rainbow/ladder \Rightarrow
conserved em current, correct charge of K^+ and K^0



charge radii	experiment	DSE calc
r_π^2	$0.44 \pm 0.01 \text{ fm}^2$	0.45 fm^2
$r_{K^+}^2$	$0.34 \pm 0.05 \text{ fm}^2$	0.38 fm^2
$r_{K^0}^2$	$-0.054 \pm 0.026 \text{ fm}^2$	-0.086 fm^2



Constituent Quark-like Behavior for c, b -quarks



- Mass shell positions marked for $\bar{b}b$ and $\bar{c}c$ quarkonia
- qQ mesons sample $M_Q(p^2) \sim 4$ times further into timelike region
- The same constituent or pole mass is unlikely to suffice for QQ and qQ mesons⁵¹



General Pseudoscalar Mass Formula

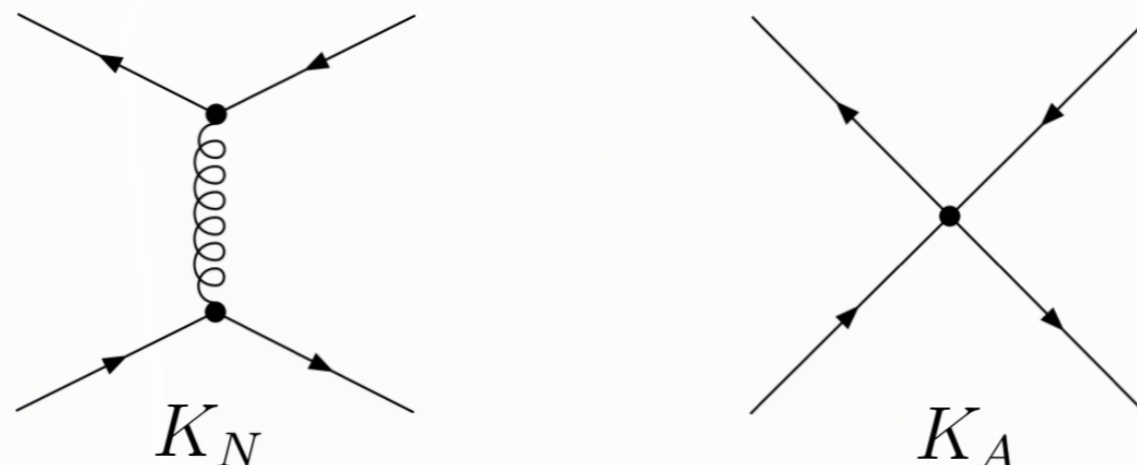
- $N_f = 3$, charge neutral states: $p = \pi^0, \eta, \eta'$

$$m_p^2 \begin{bmatrix} f_p^3 \\ f_p^8 \\ f_p^0 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ n_p \end{bmatrix} + \left[2 \mathcal{M}_{3 \times 3} \right] \begin{bmatrix} \rho_p^3 \\ \rho_p^8 \\ \rho_p^0 \end{bmatrix}$$

- Isospin breaking: $m_u \neq m_d$ allows anomaly, \mathcal{F}^0 , and $s\bar{s}$ into π^0
- η' in $SU(N_f)$ limit: $m_{\eta'}^2 f_{\eta'}^0 = n_{\eta'} + 2m \rho_{\eta'}^0$



A Schematic Model: Flavor mixing, η, η'



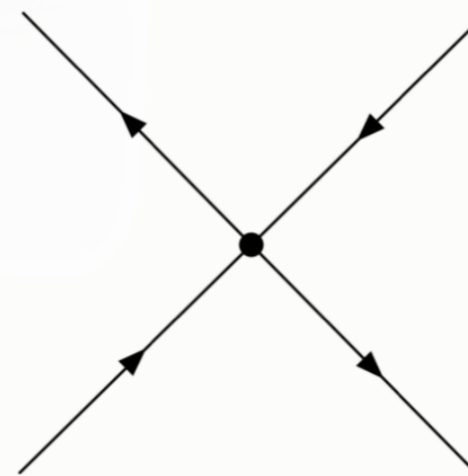
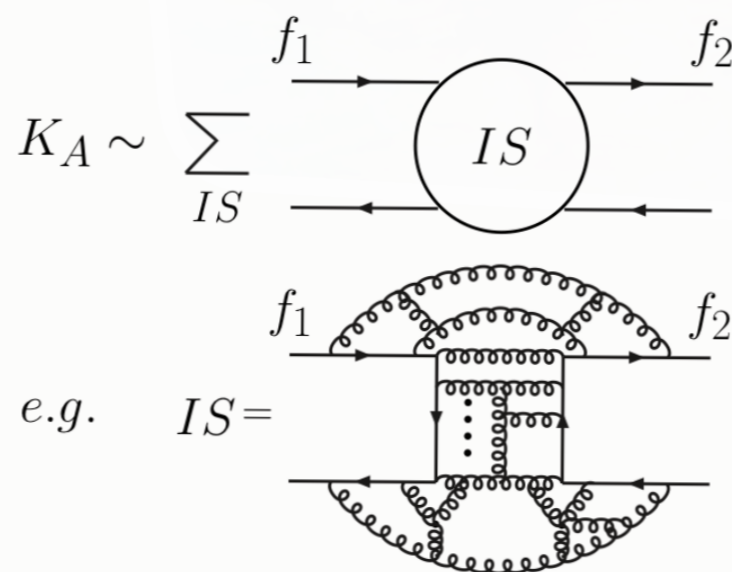
- [Bhagwat, Chang, Liu, Roberts, PCT, PRC (76), 2007; arXiv:0708.1118]
- Structure: $K_N = \text{LR vector gluon exch}$,
 $K_A = \mathcal{F}(\gamma_5, \not{P}\gamma_5) \otimes (\gamma_5, \not{P}\gamma_5)\mathcal{F}$, $\mathcal{F} = \text{diag}(1/M_f)$
- (Munczek-Nemirovsky) t-channel $\delta^4(k)$ for K_N and K_A
- 2 strength parameters: $\rho^0 \Rightarrow K_N$, $m_{\eta'} \Rightarrow K_A$.
- Fix $m_u, m_d, m_s \dots$ via vector mesons



Model Bethe-Salpeter Kernel for flavor singlet?

- Vertex integral eqns do not involve $Q_t(x)$ explicitly:

$$\Gamma_{5\mu}^\alpha(k; P) = Z_2 \gamma_5 \gamma_\mu \mathcal{F}^\alpha + \int^\Lambda K S_+ \Gamma_{5\mu}^\alpha S_-$$
- DSE models need: $K_{\text{BSE}} = K_N + K_A$, both are $\bar{q}q$ irreducible, K_N is also n-gluon irreducible

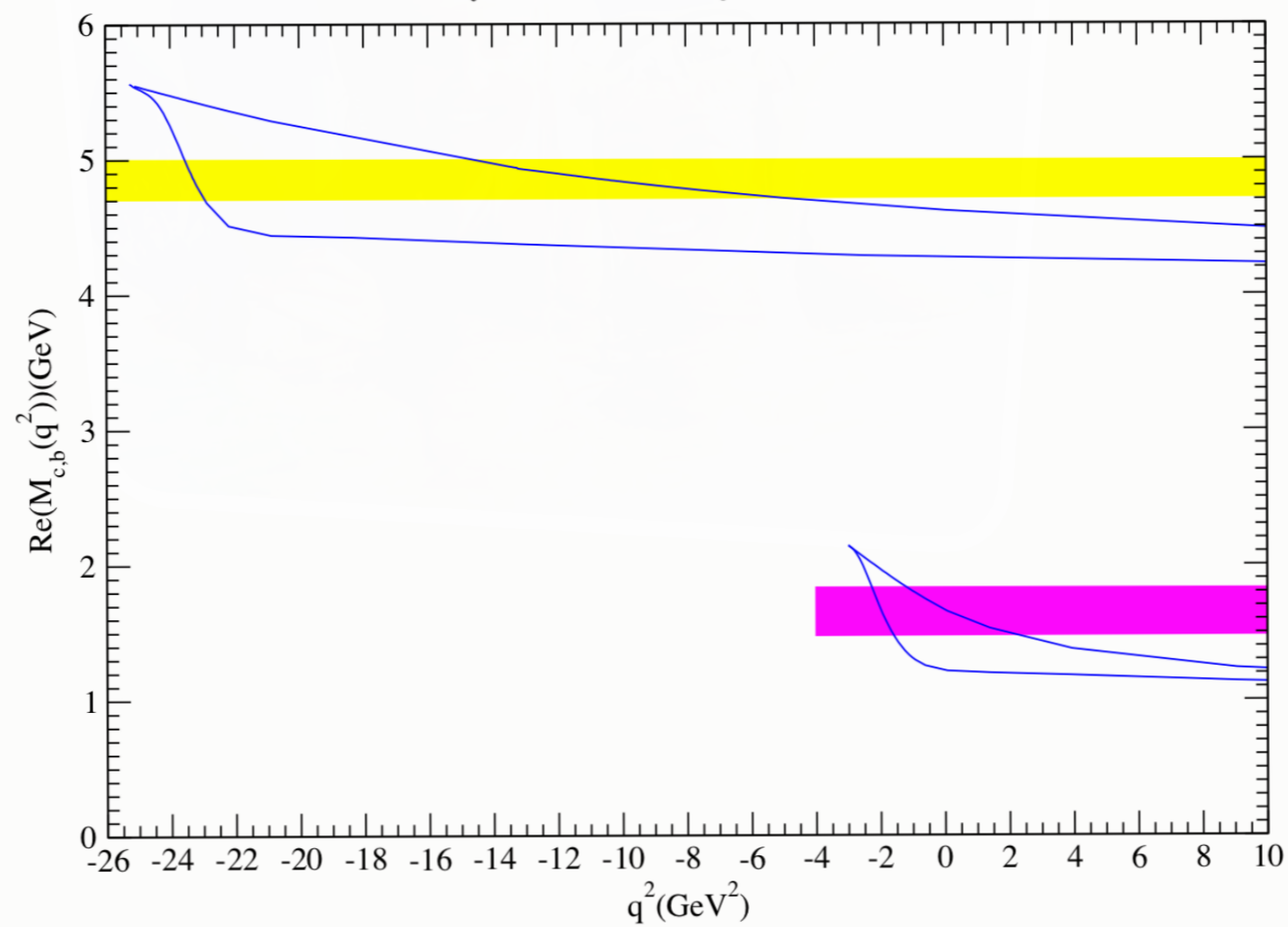


- A scenario that works: Witten-Veneziano massless axial-vector ghost linking pseudoscalar GBs



c- and b-Quark Mass Function for BSE

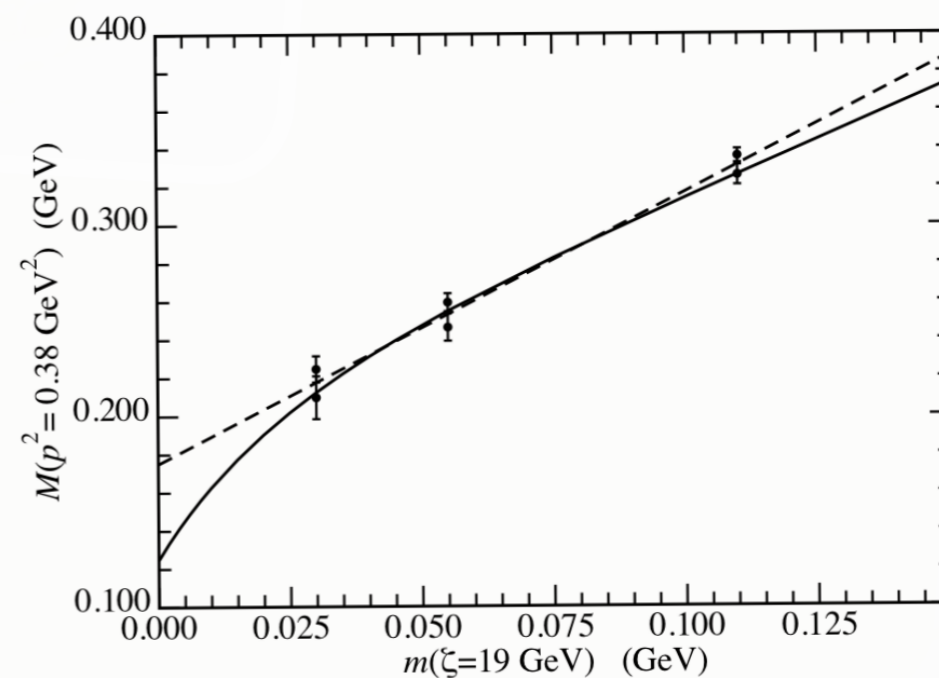
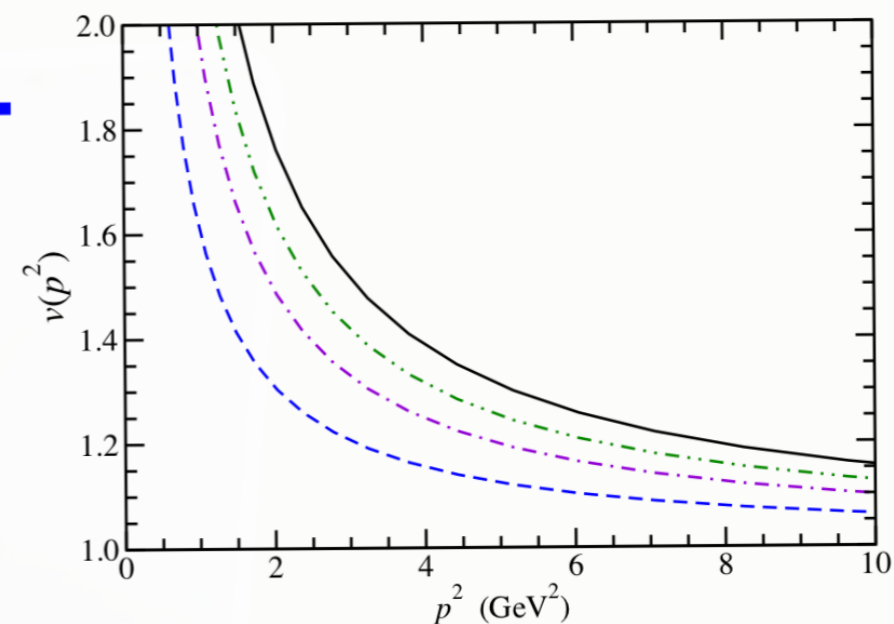
c,b quark mass function near the peak of the parabolic region with P^2 near the meson mass shells
 $m_c(19 \text{ GeV})=0.88 \text{ GeV}$, $m_b(19 \text{ GeV})=3.8 \text{ GeV}$





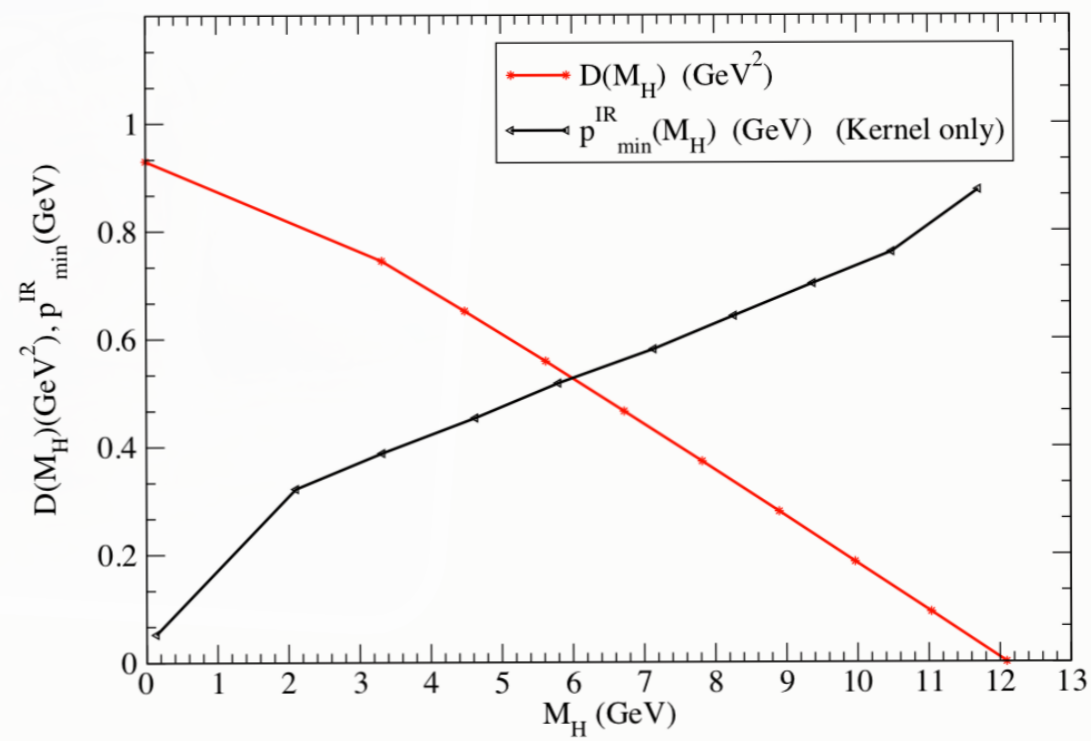
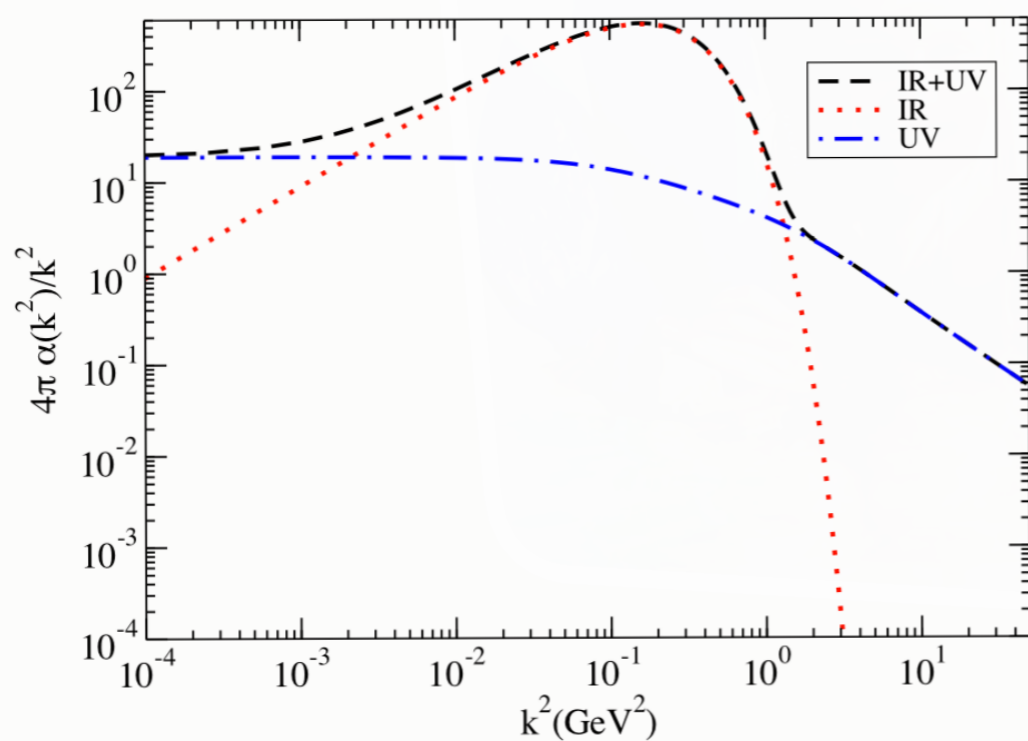
Lattice-assisted DSE Results

- Evident vertex enhancement
- Curvature in low m_q depn
- $M^{\text{IR}}(p^2)$ 40% below linear
- Chiral Extrapolation
- $\langle \bar{q}q \rangle_{\mu=1 \text{ GeV}}^{\text{qu-lat}} = -(190 \text{ MeV})^3$
- $\langle \bar{q}q \rangle^{\text{qu-lat}} \approx \langle \bar{q}q \rangle^{\text{expt}} / 2$
- f_π 30% low





IR Suppression of Kernel





Deep Inelastic Lepton Scattering

Convenient basis in Bj lim:

$$n^\nu = \frac{M}{2\omega}(1, -1; \vec{0}_\perp); \quad n^2 = 0 = p^2; \quad p \cdot n = 2 \quad .; \quad \omega = M/2 \text{ (rest frame)}, \quad \omega = \infty \text{ (IMF)}$$

$$P^\mu = \frac{M}{2}(n^\mu + p^\mu); \quad q^\mu \rightarrow \nu n^\mu + \frac{Mx}{2}(n^\mu - p^\mu) + \mathcal{O}\left(\frac{1}{\nu}\right)$$

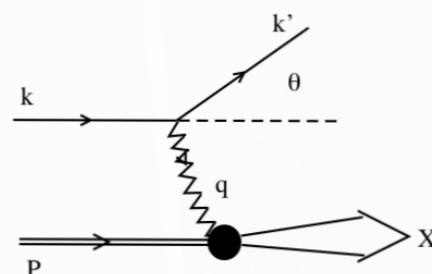
$$W^{\alpha\beta} \rightarrow (a\nu + b)(F_2 - 2x F_1) + \left(-g^{\alpha\beta} + n^\alpha \frac{P^\beta}{M} + \frac{P^\alpha}{M} n^\beta\right) F_1 + \mathcal{O}\left(\frac{1}{\nu}\right)$$

$$\{W^{\alpha\beta} q_\beta\}_{LO} = 0 = W^{\alpha\beta} n_\beta$$

handbag diagram $\Rightarrow W_{HB}^{\alpha\beta} n_\beta = 0$, (LO current consv)



Deep Inelastic Lepton Scattering



Bjorken limit:

$$\nu = q \cdot P/M \rightarrow \infty ; \quad -q^2 = Q^2 \rightarrow \infty$$

$$0 < x = \frac{Q^2}{2P \cdot q} < 1$$

$$W^{\alpha\beta} = \left\| \left[\begin{array}{c} \text{Zigzag } q \\ \text{P} \end{array} \right] \right\|^2 \sim \text{Im} \left[\left[\begin{array}{c} \text{Zigzag } q \\ \text{Zigzag } q \\ \text{P} \end{array} \right] \right] = \frac{1}{2\pi} \text{Disc } T^{\alpha\beta}(\nu)$$

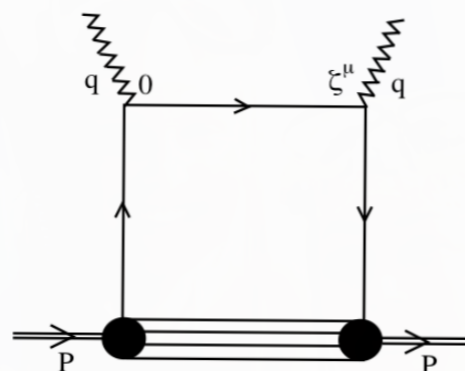
$$W^{\alpha\beta} = -\left(g^{\alpha\beta} - \frac{q^\alpha q^\beta}{q^2}\right) F_1 + \frac{P_T^\alpha(q) P_T^\beta(q)}{P \cdot q} F_2$$

$$F_1(x) = \sum_q \frac{e_q^2}{2} f_q(x) + \dots$$



Deep Inelastic Lepton Scattering

$$T^{\mu\nu}(\text{LO}) = T_{GHB}^{\mu\nu} =$$



$$q^+ = q \cdot n = -Mx, \quad |\xi^-| \sim \frac{1}{Mx}$$

$$q^- = q \cdot p = 2\nu, \quad |\xi^+| \sim 0$$

DIS is hard and fast—confinement is soft and slow $\Rightarrow S(k+q) \rightarrow \frac{\gamma^+}{2(k^+ - P^+ x) + i\epsilon}$

$W^{\mu\nu} \propto \{T^{\mu\nu}(\epsilon) - T^{\mu\nu}(-\epsilon)\} \Rightarrow$ Euclidean model elements can be continued

EG, π^+ target : $f_q(x) = \frac{1}{4\pi} \int d\xi^- e^{iq^+\xi^-} \langle \pi(P) | \bar{q}(\xi^-) \gamma^+ q(0) | \pi(P) \rangle_c = -f_{\bar{q}}(-x)$

$$f_q(x) = \frac{1}{2} \text{tr} \int \frac{d^4k}{(2\pi)^4} \delta(k^+ - P^+ x) S(k) \gamma^+ S(k) T(k, P)$$

General $T(k, P) = \bar{u}\pi^+$ scattering amplitude:

s-channel structure \rightarrow "spectator \bar{d} " $\Rightarrow f_u(x), \quad 0 < x < 1$

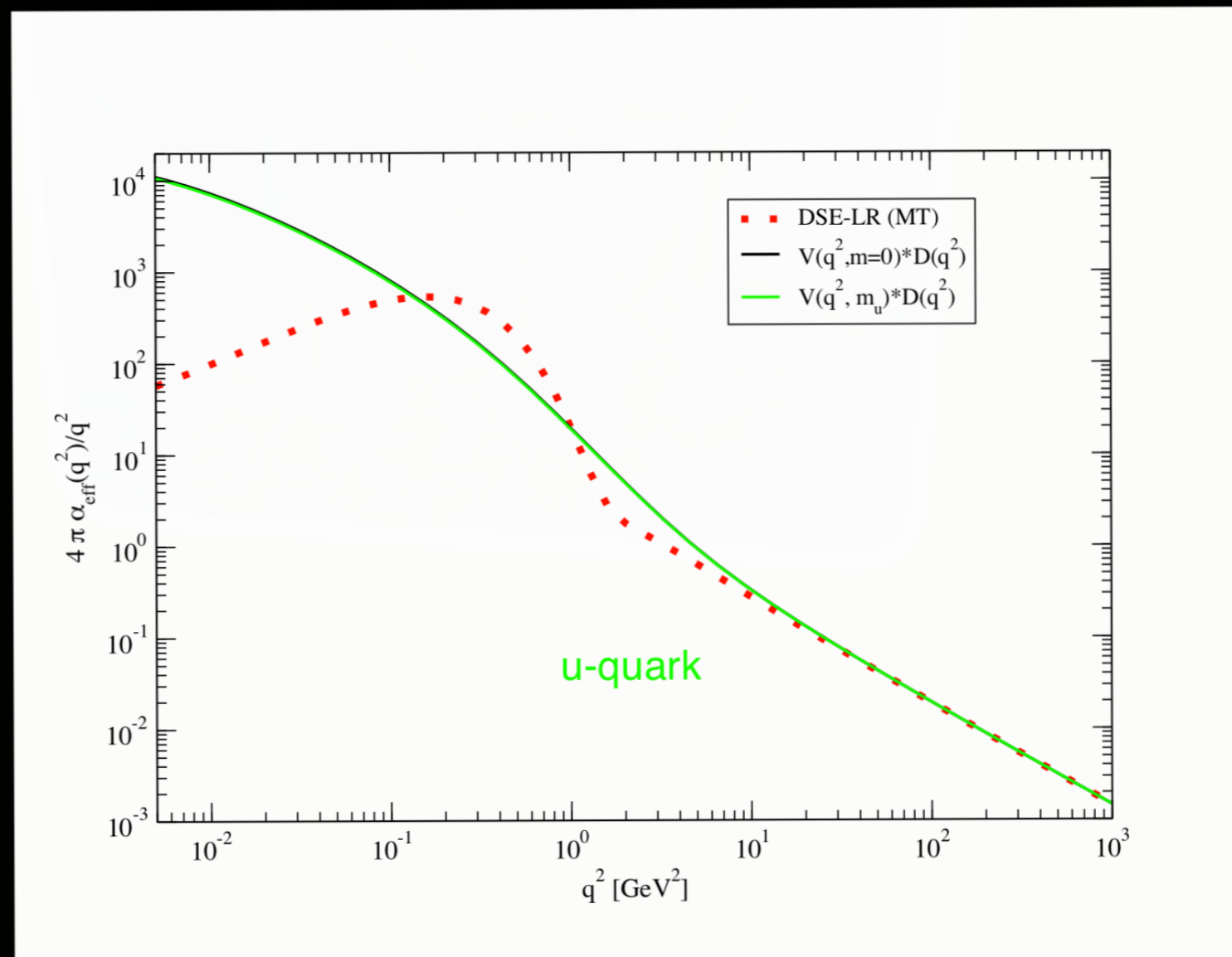
u-channel structure \rightarrow "spectator $u\bar{d}$ " $\Rightarrow f_{\bar{u}}(-x), \quad 0 < x < 1$

correct x
support





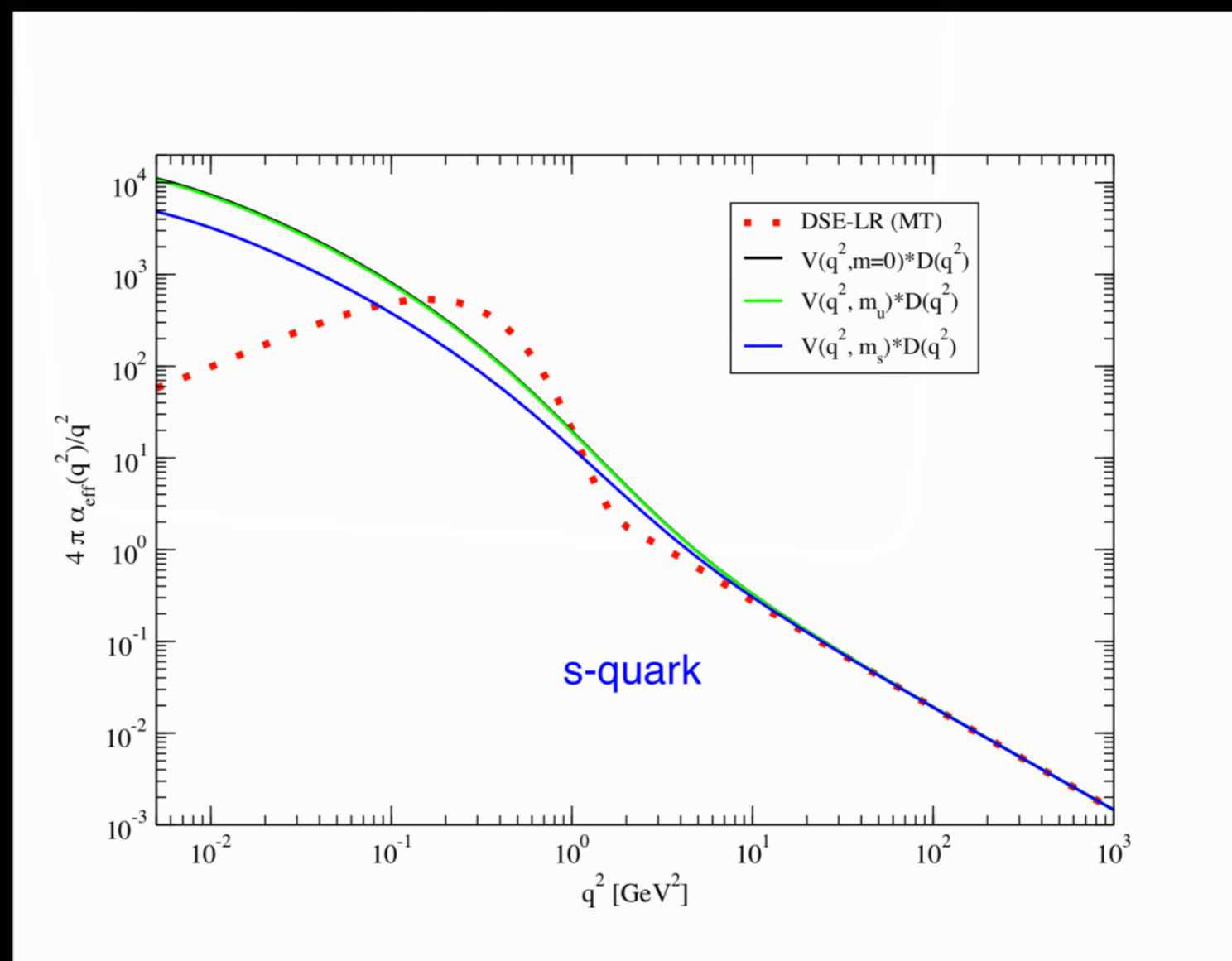
Quenched lattice $\Rightarrow m_q$ Depn of DSE Kernel



Bhagwat, Pichowsky, Roberts, Tandy, PRC68, 015203 (2003)



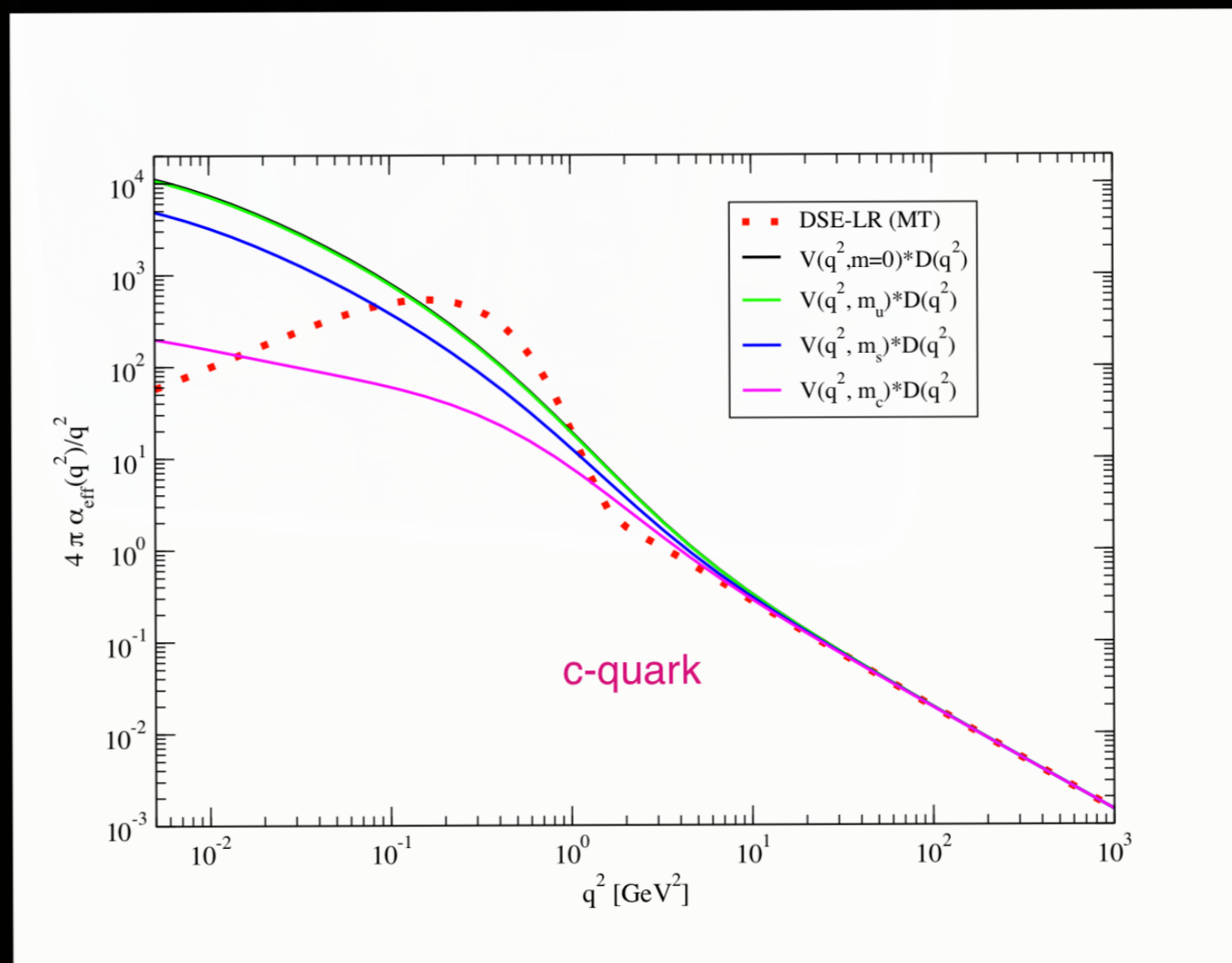
Quenched lattice $\Rightarrow m_q$ Depn of DSE Kernel



Bhagwat, Pichowsky, Roberts, Tandy, PRC68, 015203 (2003)



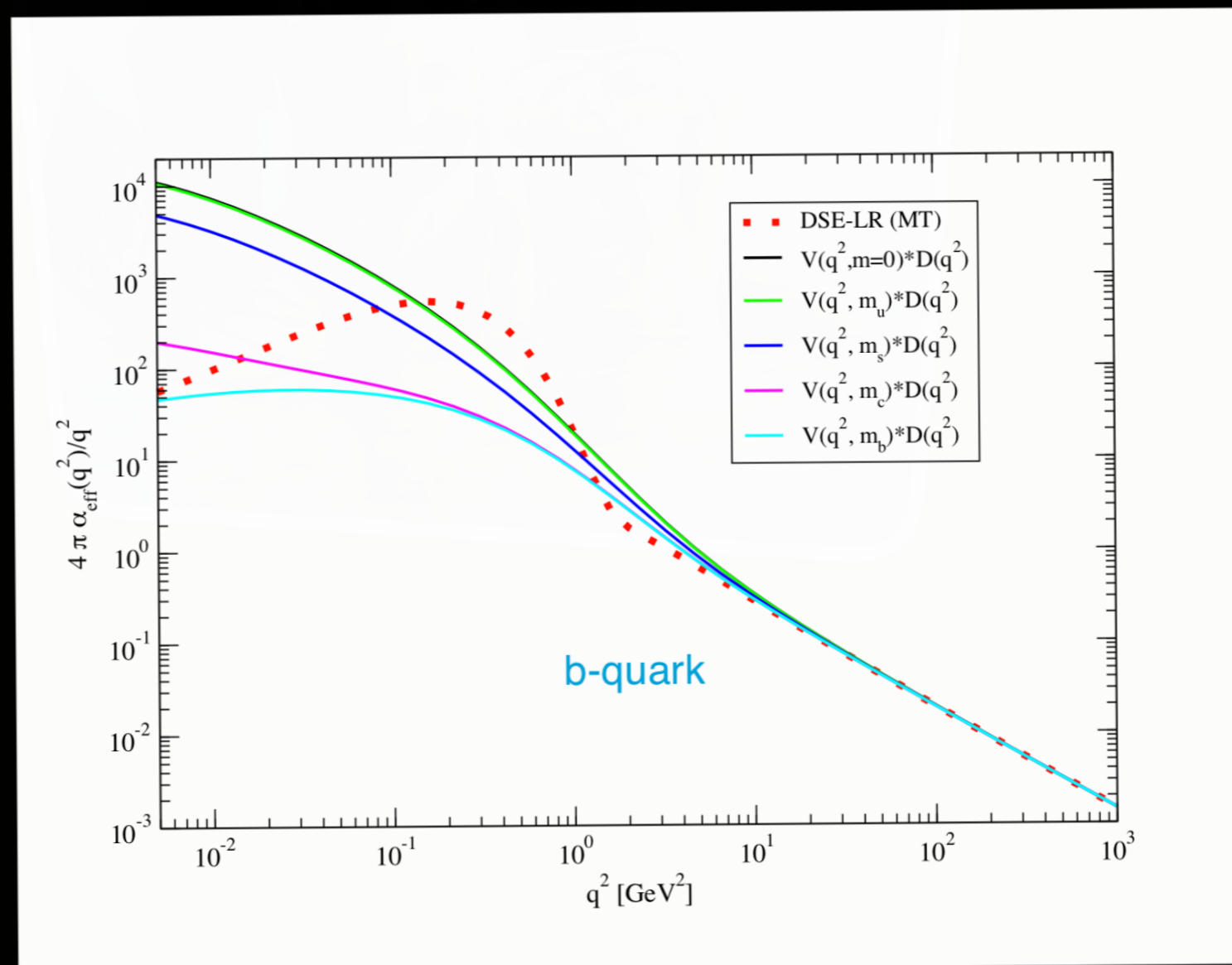
Quenched lattice $\Rightarrow m_q$ Depn of DSE Kernel



Bhagwat, Pichowsky, Roberts, Tandy, PRC68, 015203 (2003)



Quenched lattice $\Rightarrow m_q$ Depn of DSE Kernel



Bhagwat, Pichowsky, Roberts, Tandy, PRC68, 015203 (2003)